Probing many-body localized phase and delocalization transition with matrix elements

> Maksym Serbyn UC Berkeley



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### Artificial quantum systems

mutli-qubit systems



#### cold atoms



trapped ions



# NV centers in diamond, polar molecules,



Universality in **isolated quantum systems** out-of-equilibrium?

#### **Outcomes of unitary dynamics**

• Thermalization:

$$\begin{array}{c} & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ &$$

and the second second

MBL phase: breakdown of thermalization

[Anderson, Fleishman'80] [Basko, Aleiner, Altshuler'05] [Gorniy, Polyakov, Mirlin'05] [Oganesyan, Huse'08] [Znidaric, Prosen'08] [Pal, Huse'10]

#### **Experiments on MBL**





• MBL signatures in dynamics? Local observables?



#### **Quasi-local conserved quantities**

• If model is in MBL phase, rotate basis

• Conserved spins:  $\tau_i = U^{\dagger} S_i U$  quasi-local; complete basis



 Consequences: no transport, ETH breakdown, power-law relaxation

[MS, Papic, Abanin '13] [Huse, Oganesyan '13] [Imbrie'14]

I.

#### **Dynamics in MBL phase**



- Logarithmic growth of entanglement [MS, Papic, Abanin, PRL'13]
- Power-law relaxation of local observables:

$$\left|\langle\widehat{\mathcal{O}}(t)
ight
angle - \langle\mathcal{O}(\infty)
ight
angle \right| \sim rac{1}{t^a}$$
 memory of initial state

[MS, Papic, Abanin, PRB'14]

#### Outline



- 1. Delocalization transition [MS, Papic, Abanin, PRX'15]
- 2. Entanglement spectrum in MBL phase

[MS, Michailidis, Abanin, Papic, arXiv:1605.05737]





Alternatives? Scaling parameter for transition?
 [MS, Papic, Abanin, PRX'15]

#### Thouless conductance in Anderson localization

How to quantify what happens if we join two systems?

- Coupling **vs** Level spacing
  - # g>>1: states become extended
  - \* g<<1: localization!
- Scaling theory of localization:

$$g = \frac{\sigma L^{d-2}}{e^2/h}$$

[Abrahams,Anderson, Licciardello,Ramakrishnan]

Practical definition of g:

$$g = \frac{1}{\Delta} \frac{\partial^2 E}{\partial \phi^2}$$





#### Matrix elements of local operators

$$\begin{array}{c} \text{local} \\ \text{perturbation} \\ \downarrow \\ \uparrow \\ \hline \\ R \end{array}$$

Matrix elements from ETH

$$\langle i|S^z|j\rangle = e^{-S(E,R)/2}f(E_i,E_j)R_{ij}$$

[Srednicki'99]

narrow distribution:  $\langle i|S^z|j
angle \sim 1/\sqrt{2^R}$ 

• Matrix elements from  $\tau_i$ 

$$S^{z} = \sum_{\{\alpha\}} \hat{\tau}^{\{\alpha\}} \hat{B}^{\{\alpha\}} [\tau^{z}]$$
$$\langle i|S^{z}|j \rangle$$

broad distribution:  $\langle i|S^{z}|j\rangle \sim \exp(-\kappa' R)$ 



#### Many-body analogue of Thouless conductance

- Effect of local perturbation on eigenstates:  $H \rightarrow H + V$
- $H|n\rangle = E_n|n\rangle$   $(H+V)|\alpha\rangle = E_\alpha|\alpha\rangle$
- New eigenstates are localized/delocalized?
- Parameter:

 $\mathcal{G} = \log \frac{V_{i,i+1}}{E_i - E_{i+1}}$ 



 $V_{nm}$ 

 $\mathcal{G}\gg 1$ 

strong mixing all spins perturbed





 $\mathcal{G}\ll 1$ no resonances au are local

#### **Distribution of Thouless conductance**

- Scaling parameter:  $\mathcal{G}(L) = \log \frac{V_{n,n+1}}{|E_n E_{n+1}|}$
- Numerical results for XXZ spin chain: [MS, Papic, Abanin, PRX'15]

 $\widehat{V} = S_1^z \quad \bigstar \quad \uparrow^{J_\perp} \\ J_z \downarrow \quad \uparrow^{h_i}$ disorder W Thermalizing phase **MBL** phase  $5 \frac{x 10 \bar{x}^3 10 \bar{x}^3}{5 \bar{x}^3 \bar{x}^3} W = 5$  $7 \frac{x 10x^{-3}10x^{-3}}{10x^{-3}10^{-3}}$ W = 3.6W = 0.50.0105.<del>0105.015</del> 8 8 <del>0 10</del> 10 <del>0 10</del> 10 0-10-10  $\frac{2 - 12}{12}$  12 2-12-12 2-12-12 0.010.010.0114-14-14 **4-14**-14 4-14-14 3 <del>6-16-</del>16 <del>316-16-</del>16 <del>416-16-</del>16  $p(\mathcal{G})$  $p(\mathcal{G})$ 3 3 2 2 0.005.005.005 2 | 2|2 1 1 -6 -6 -46 -4 -24 -2 02 0 20 2 42 4 64 6 6-15-15-15-150-10-10-10-5-50 0 05 5 5 0 6 G G G

#### Scaling of G and many-body mobility edge

- Energy resolved  $\mathcal{G}$ :  $\mathcal{G}(L,\varepsilon) = \left\langle \log \frac{\langle i|S^z|j\rangle}{|E_i E_j|} \right\rangle$
- Numerical results for XXZ spin chain
- Exponent  $v=0.7\pm0.1$  agrees with numerics for L=22 spins

[Luitz, Laflorencie, Alet PRB'15]





#### Dynamics at the transition

- Qualitative argument: G(L) = const+O(log L) at transition  $\rightarrow$ logarithmically slow transport [MS, Papic, Abanin, PRX'15]
- tEBD simulations for L=24 spins:



$$\langle S_{z,L}^2 \rangle - \langle S_{z,L} \rangle^2 \propto \log t$$

$$= \frac{W=2.5}{W=3}$$

$$W=3.5$$

$$W=5$$

v 2



Consistent with RG studies

[Vosk,Huse,Altman,PRX'15] [Potter,Vasseur,Parameswaran,PRX'15]

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#### From entanglement entropy to spectrum

• "Quantumness" of the pure state:



trace out  $R \rightarrow$  $\rho_L = \operatorname{Tr}_R |\psi\rangle\langle\psi|$ 

- Entanglement entropy:  $S_{\text{ent}} = -\sum_i \lambda_i \log \lambda_i$ 
  - # ground states: probes topological order
    [Levin&Wen], [Kitaev&Preskill]
  - \* excited states: probes ergodicity
- Beyond entanglement? More information in  $\{\lambda_i\}$ [Li & Haldane]

#### Organization of entanglement spectrum



MBL phase: conserved quantities label ES

$$|\uparrow\uparrow\uparrow\uparrow\rangle = c_0 |\uparrow\uparrow\rangle|\uparrow\uparrow\rangle + e^{-\kappa}|\uparrow\downarrow\rangle|\uparrow\uparrow\rangle + e^{-2\kappa}|\uparrow\downarrow\rangle|\downarrow\uparrow\rangle + \dots$$
$$\stackrel{\leftarrow}{r=1} \qquad \stackrel{\leftarrow}{r=2} + e^{-4\kappa} |\downarrow\downarrow\rangle|\downarrow\downarrow\rangle + \dots$$
$$\stackrel{\leftarrow}{r=4} + \dots$$

• Coefficients decay as  $|C_{\uparrow...\uparrow\downarrow\downarrow\uparrow\uparrow\downarrow\uparrow...\uparrow}| \propto e^{-\kappa r}$ 

#### Power-law entanglement spectrum

• Hierarchical structure of  $ho_L = \sum_{r=0}^L |\psi^{(r)}\rangle \langle \psi^{(r)}|$ 

 $\langle \psi^{(r)} | \psi^{(r)} \rangle \propto e^{-2\kappa r}$ 

Orthogonalize perturbatively

 $\lambda^{(r)} \propto e^{-4\kappa r}$ 

multiplicity is  $2^{r}$ 

but non-orthogonal

Power-law entanglement spectrum

$$\lambda_k \propto rac{1}{k^\gamma} \qquad \gamma pprox rac{4\kappa}{\ln 2}$$

#### Numerics for XXZ spin chain

Numerical studies for XXZ spin chain,  $J \perp = J_z = 1$ 





Power law entanglement spectrum:



more details in: [arXiv:1605.05737]

#### Decay of entanglement spectrum

•  $\gamma$  controls decay of entanglement spectrum  $\lambda_k \propto rac{1}{k\gamma}$ 



 $\gamma \approx \frac{4\kappa}{\ln 2}$ 

perturbation theory  $\kappa = 2\kappa' + \ln 2$  $\mathcal{G}(L) \propto e^{-\kappa' L}$ 

• Large value of  $\gamma \rightarrow MPS$  description!

#### MPS algorithm close to MBL transition

• Theoretically:  $\gamma$  controls truncation error  $\propto 1/\chi^{\gamma-1}$ 

large  $\gamma \rightarrow$  MPS is ok close to MBL transition

 Practically: MPS-based algorithm in MBL phase use ES as a very stringent test for DMRG:



[arXiv:1605.05737]

also: [Yu et al arXiv:1509.01244] [Lim&Sheng arXiv:1510.08145] [Pollmann et al arXiv:1509.00483] [Kennes&Karrasch arXiv:1511.02205]

#### Estimates for the bond dimension

- Entanglement spectrum deviates at small  $\lambda$
- Average bond dimension to converge *S<sub>ent</sub>* up to 1%:



Uses: DMRG close to MBL transition, probe MBL phase

#### Summary

- Scaling parameter for transition  $\mathcal{G} = \log \frac{V_{i,i+1}}{E_i E_{i+1}}$  [PRX 5, 041047 (2015)]  $\rightarrow$  logarithmic transport at transition
- Power-law entanglement spectrum in MBL  $\lambda_k \propto \frac{1}{k\gamma}$  $\rightarrow$  power  $\gamma$  related to scaling of  $\mathscr{G}$  [arXiv:1605.05737]

→implementation of MPS algorithm close to transition

Global goal: delocalization transition & thermalization





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## GORDON AND BETTY FOUNDATION

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