

# Anderson Localization: Multifractality, Symmetries, Topologies, and Interactions

Alexander D. Mirlin  
Karlsruhe Institute of Technology

Moscow Institute of Physics and Technology, 8 October 2014

## Plan

- Anderson localization: basic properties, field theory
- Wave function multifractality
- Symmetries of disordered systems
- Manifestations of topology in localization theory
- Influence of electron-electron interaction

# Anderson localization



Philip W. Anderson

1958 “Absence of diffusion  
in certain random lattices”

sufficiently strong disorder → quantum localization

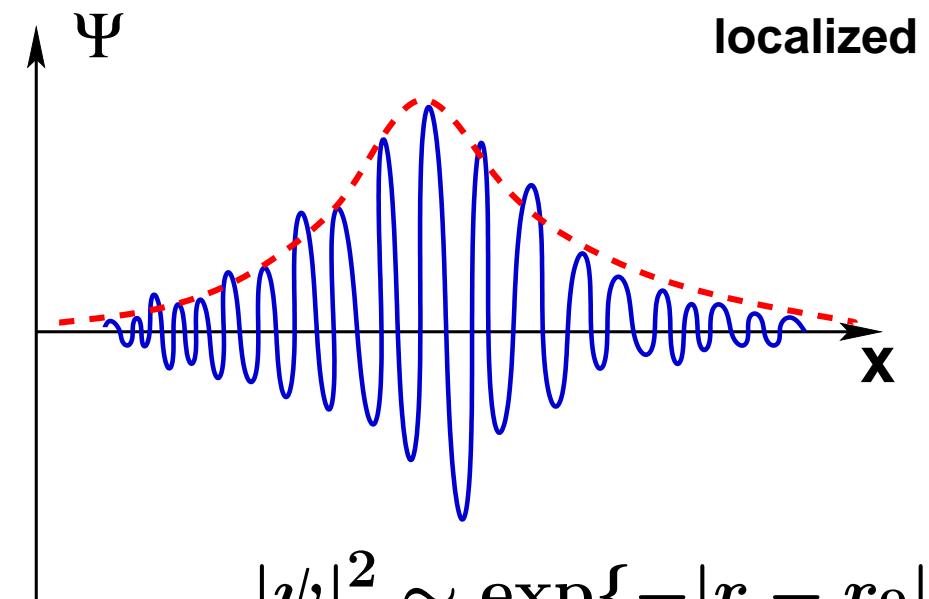
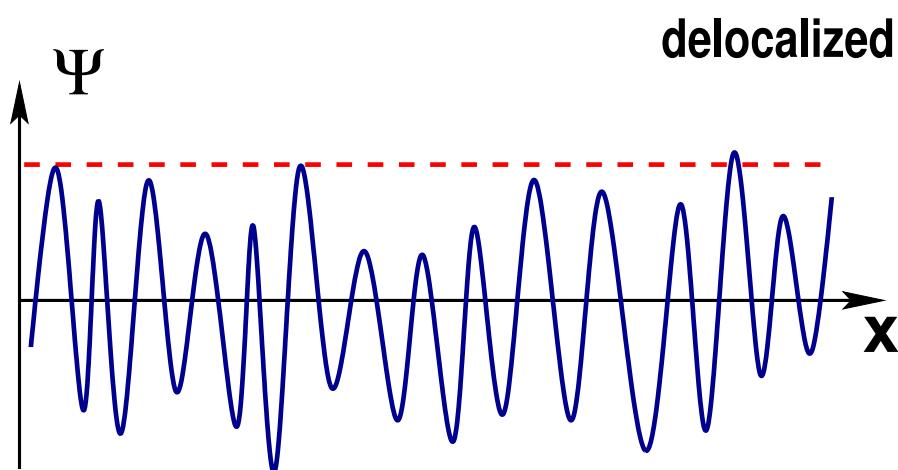
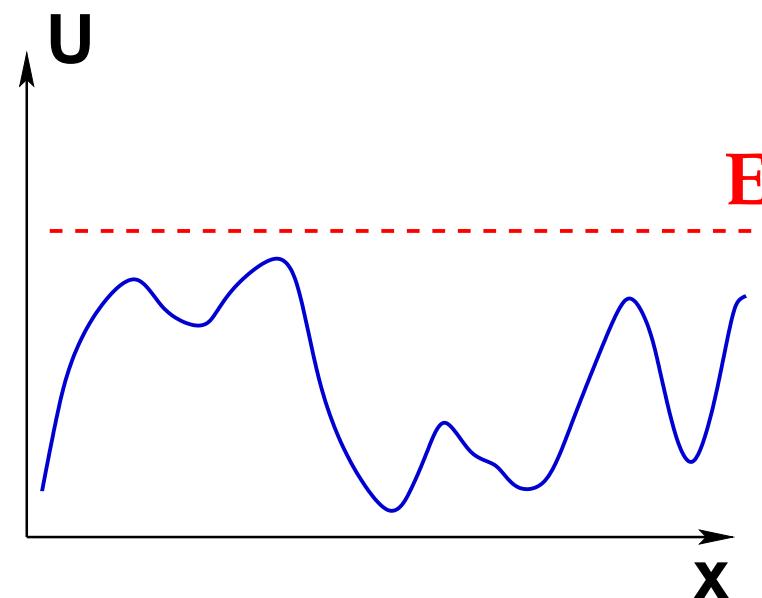
- eigenstates exponentially localized, no diffusion
- Anderson insulator

Nobel Prize 1977

# Anderson Localization: Extended and localized wave functions

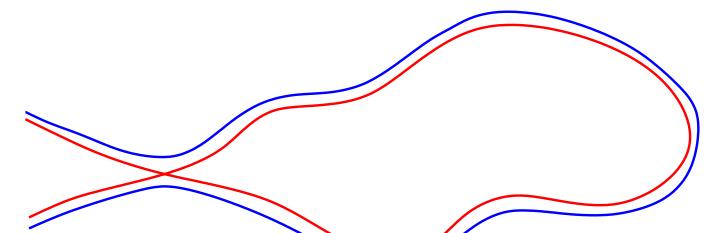
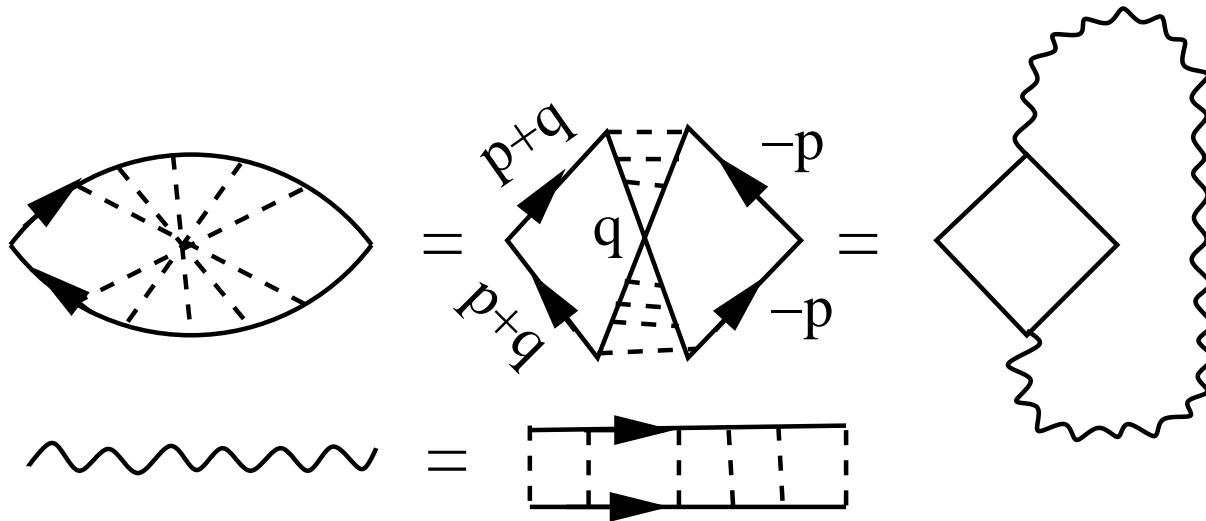
Schrödinger equation  
in a random potential

$$[-\hbar^2 \frac{\Delta}{2m} + U(r)]\psi = E\psi$$



$$|\psi|^2 \sim \exp\{-|r - r_0|/\xi\}$$

# Precursor of strong Anderson localization: Weak localization



Cooperon loop (interference of time-reversed paths)

$$\Delta\sigma_{WL} \simeq \sigma_0 \frac{1}{\pi\nu} \int \frac{(dq)}{Dq^2 - i\omega}$$

$$\Delta\sigma_{WL} = -\frac{e^2}{2\pi h} \left( \frac{\sim 1}{l} - \frac{1}{L_\omega} \right), \quad \text{3D}$$

$$L_\omega \simeq (D/\omega)^{1/2}$$

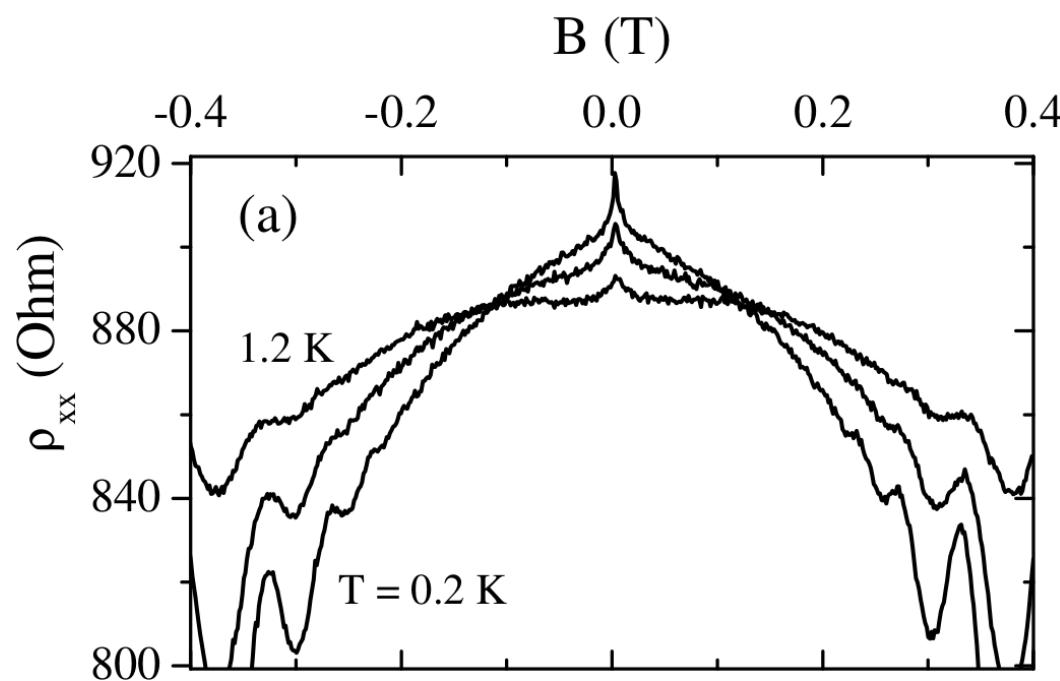
$$\Delta\sigma_{WL} = -\frac{e^2}{\pi h} \ln \frac{L_\omega}{l}, \quad \text{2D}$$

$$\Delta\sigma_{WL} = -\frac{e^2}{h} L_\omega, \quad \text{quasi-1D}$$

Generally: IR cutoff  
 $L_\omega \rightarrow \min\{L_\omega, L_\phi, L, L_H\}$

$e^2/h \simeq (25 \text{ } k\Omega)^{-1}$   
 “conductance quantum”

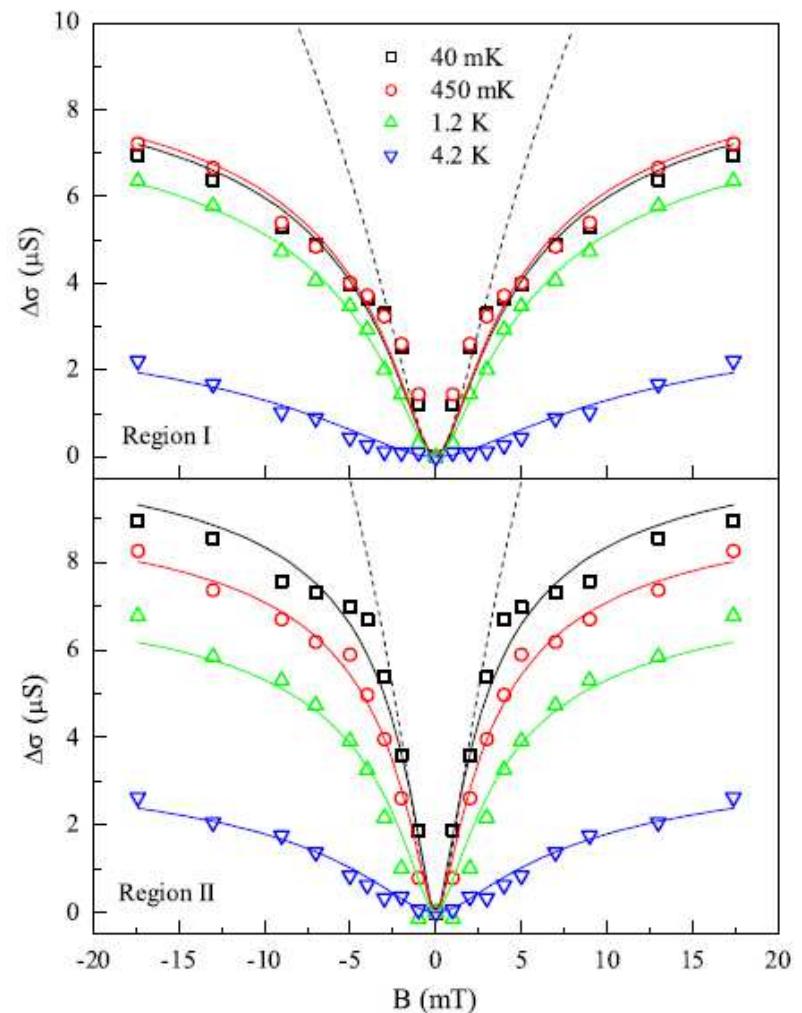
# Weak localization in experiment: Magnetoresistance



Li et al. (Savchenko group), PRL'03

2D electron gas  
in GaAs heterostructure

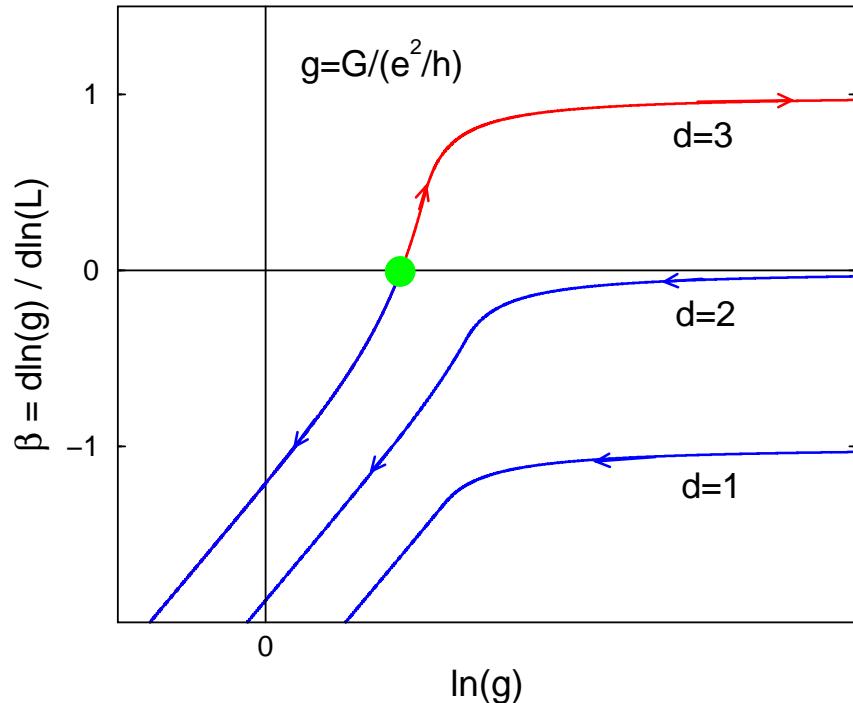
low field: weak localization



Gorbachev et al. (Savchenko group),  
PRL'07

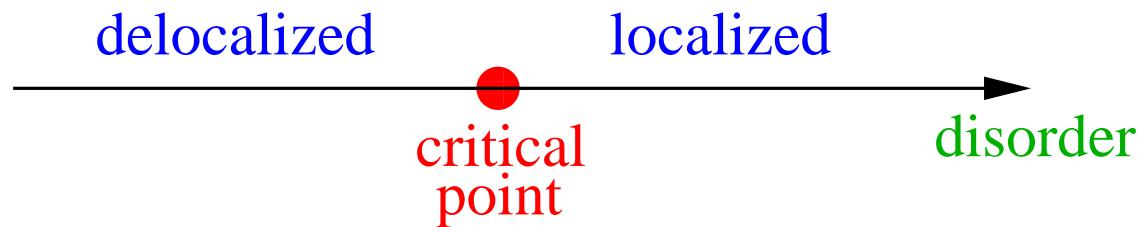
weak localization in bilayer graphene

# Anderson Insulators & Metals



quasi-1D, 2D :  
all states are localized

$d > 2$ : Anderson metal-insulator transition



review: Evers, ADM, Rev. Mod. Phys. 80, 1355 (2008)

Connection with scaling theory of critical phenomena: Thouless '74; Wegner '76

Scaling theory of localization:  
Abrahams, Anderson, Licciardello,  
Ramakrishnan '79

scaling variable:  
dimensionless conductance  $g = G/(e^2/h)$

RG for field theory ( $\sigma$ -model)  
Wegner '79

## Field theory: non-linear $\sigma$ -model

action:

$$S[Q] = \frac{\pi\nu}{4} \int d^d r \operatorname{Tr} [-D(\nabla Q)^2 - 2i\omega\Lambda Q], \quad Q^2(r) = 1$$

Wegner'79

$\sigma$ -model manifold:

e.g., “unitary” symmetry class (broken time-reversal symmetry):

- fermionic replicas:  $U(2n)/U(n) \times U(n)$  ,  $n \rightarrow 0$       “sphere”
- bosonic replicas:  $U(n,n)/U(n) \times U(n)$  ,  $n \rightarrow 0$       “hyperboloid”
- supersymmetry (Efetov'83):  $U(1,1|2)/U(1|1) \times U(1|1)$
- {“sphere”  $\times$  “hyperboloid”} “dressed” by anticommuting variables
- with electron-electron interaction: Finkelstein'83

## $\sigma$ model: Perturbative treatment

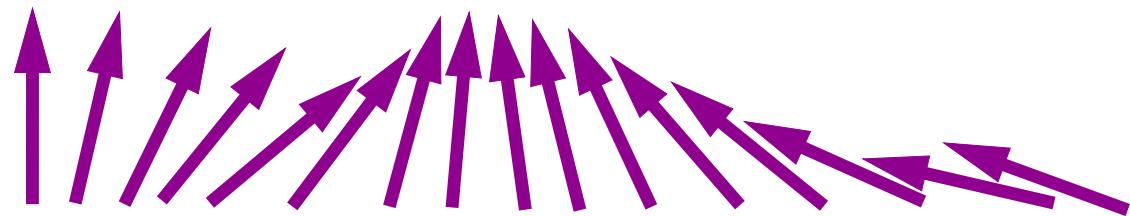
For comparison, consider ferromagnet model in external magnetic field:

$$H[S] = \int d^d r \left[ \frac{\kappa}{2} (\nabla S(r))^2 - BS(r) \right], \quad S^2(r) = 1$$

$n$ -component vector  $\sigma$ -model

Target manifold:

sphere  $S^{n-1} = O(n)/O(n-1)$



Independent degrees of freedom: transverse part  $S_\perp$ ;  $S_1 = (1 - S_\perp^2)^{1/2}$

$$H[S_\perp] = \frac{1}{2} \int d^d r \left[ \kappa [\nabla S_\perp(r)]^2 + BS_\perp^2(r) + O(S_\perp^4(r)) \right]$$

Ferromagnetic phase: broken symmetry,

Goldstone modes – spin waves

$$\langle S_\perp S_\perp \rangle_q \propto \frac{1}{\kappa q^2 + B}$$

Similarly

$$S[Q] = \frac{\pi\nu}{4} \int d^d r \text{Str}[D(\nabla Q_\perp)^2 - i\omega Q_\perp^2 + O(Q_\perp^3)]$$

theory of “interacting” diffusion modes;

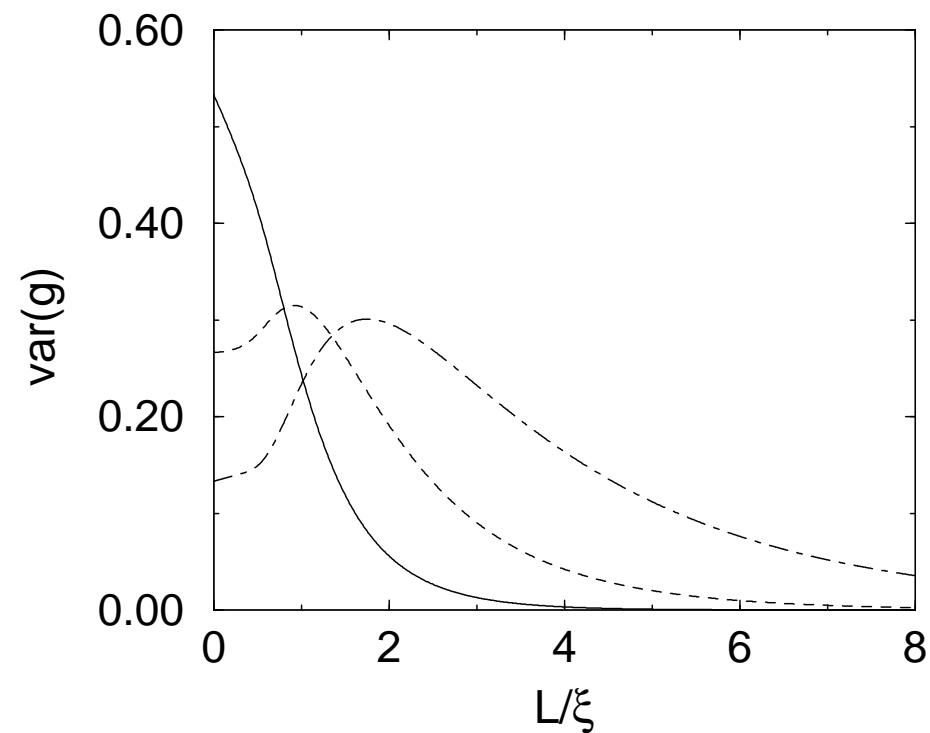
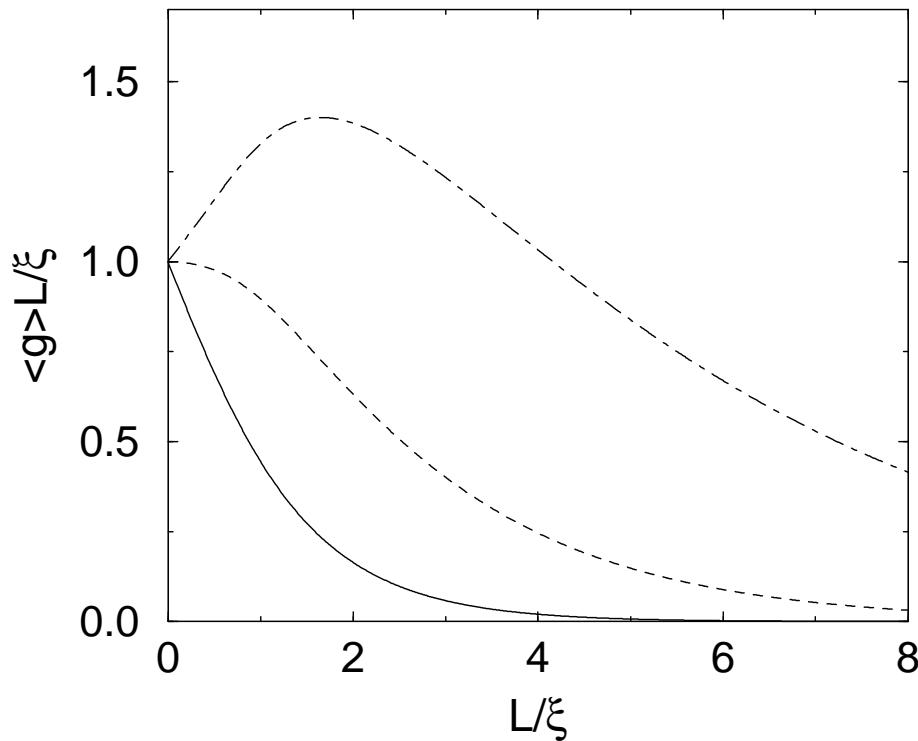
Goldstone mode: diffusion propagator

$$\langle Q_\perp Q_\perp \rangle_{q,\omega} \sim \frac{1}{\pi\nu(Dq^2 - i\omega)}$$

# Quasi-1D geometry: Exact solution of the $\sigma$ -model

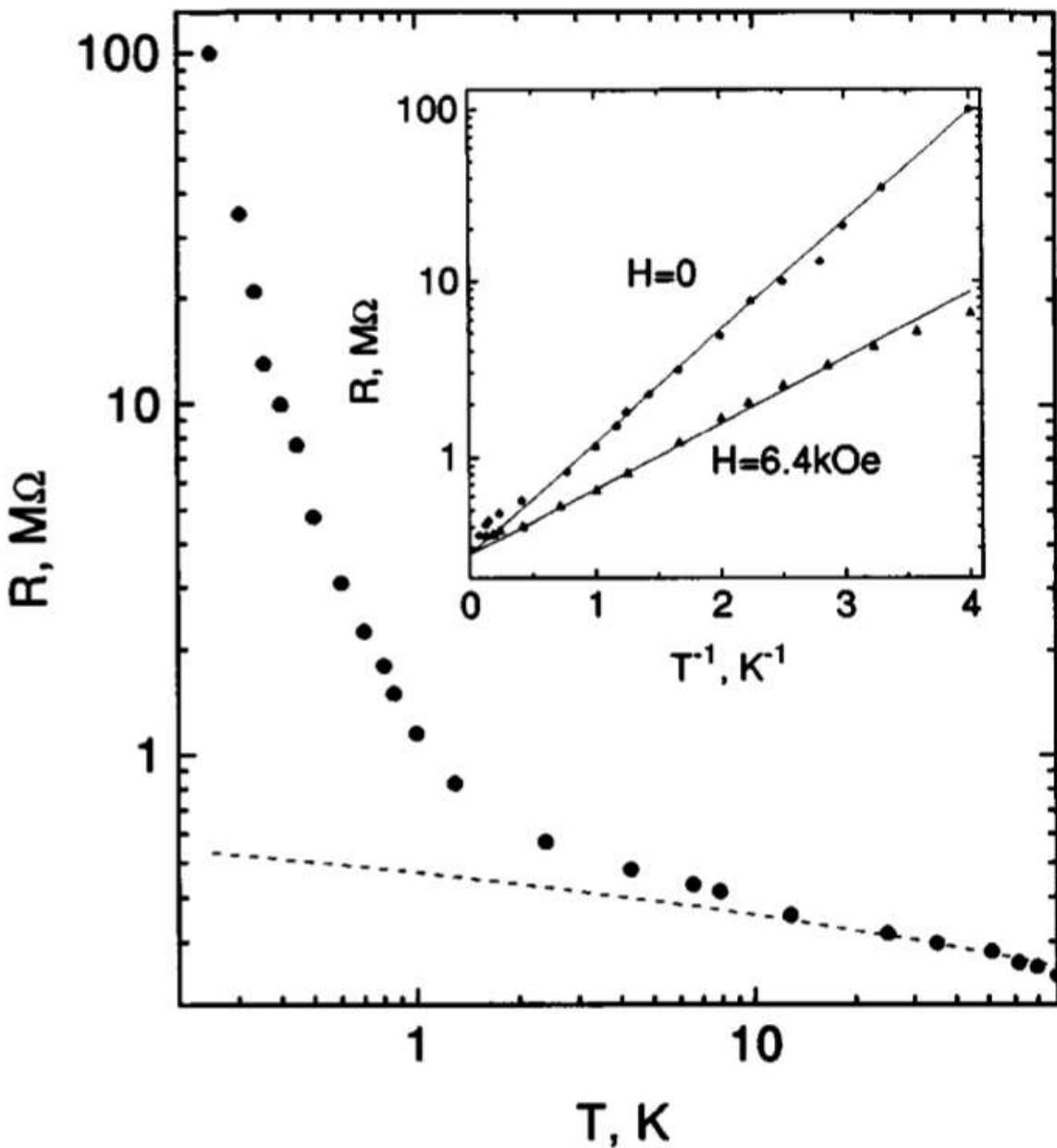
quasi-1D geometry (many-channel wire)  $\longrightarrow$  1D  $\sigma$ -model  
 $\longrightarrow$  diffusion on  $\sigma$ -model curved space  $\partial_t W = \Delta_Q W$ ,  $t = x/\xi$

- Localization length Efetov, Larkin '83
- Exact solution for the statistics of eigenfunctions Fyodorov, ADM '92-94
- Exact  $\langle g \rangle(L/\xi)$  and  $\text{var}(g)(L/\xi)$  Zirnbauer, ADM, Müller-Groeling '92-94



orthogonal (full), unitary (dashed), symplectic (dot-dashed)

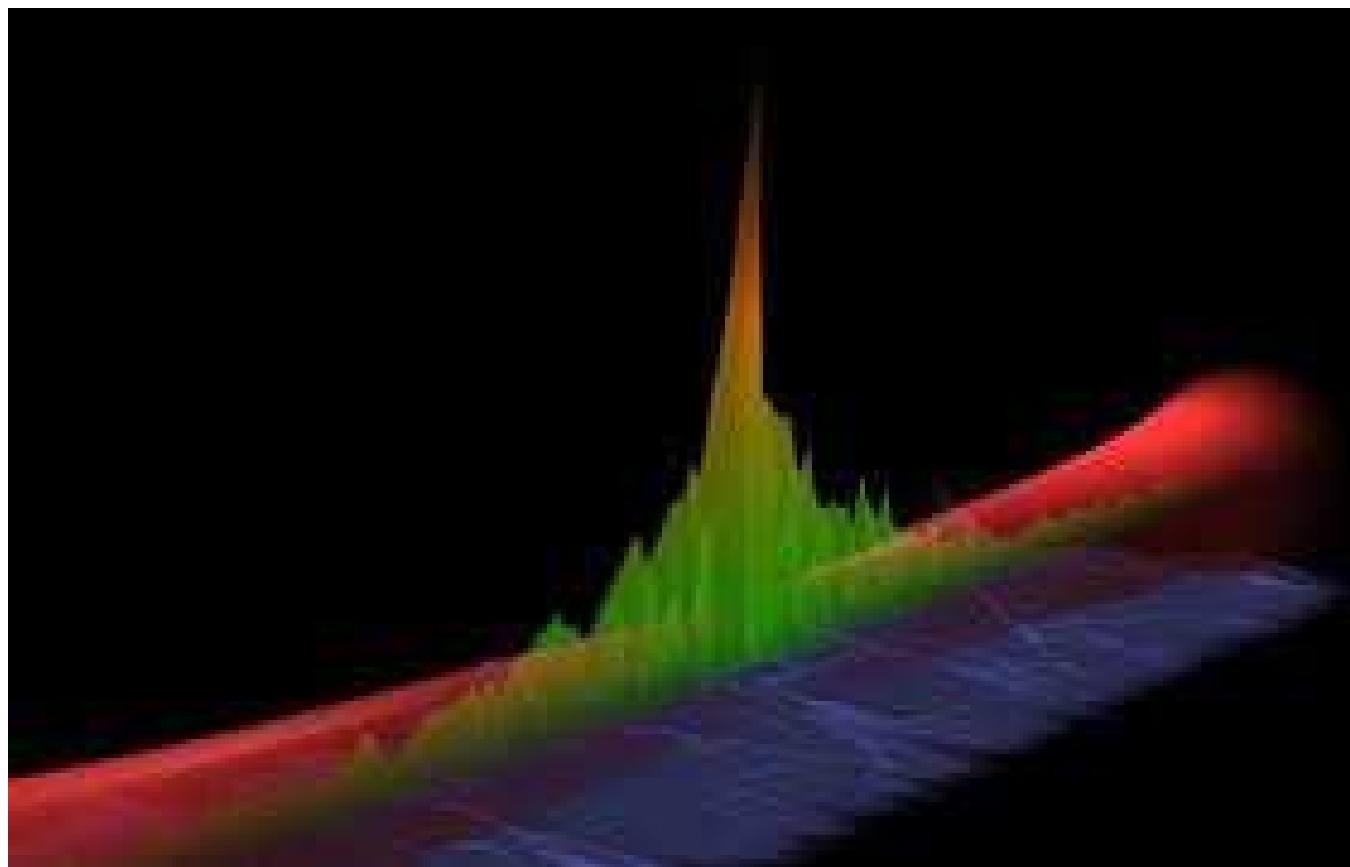
# From weak to strong localization of electrons in wires



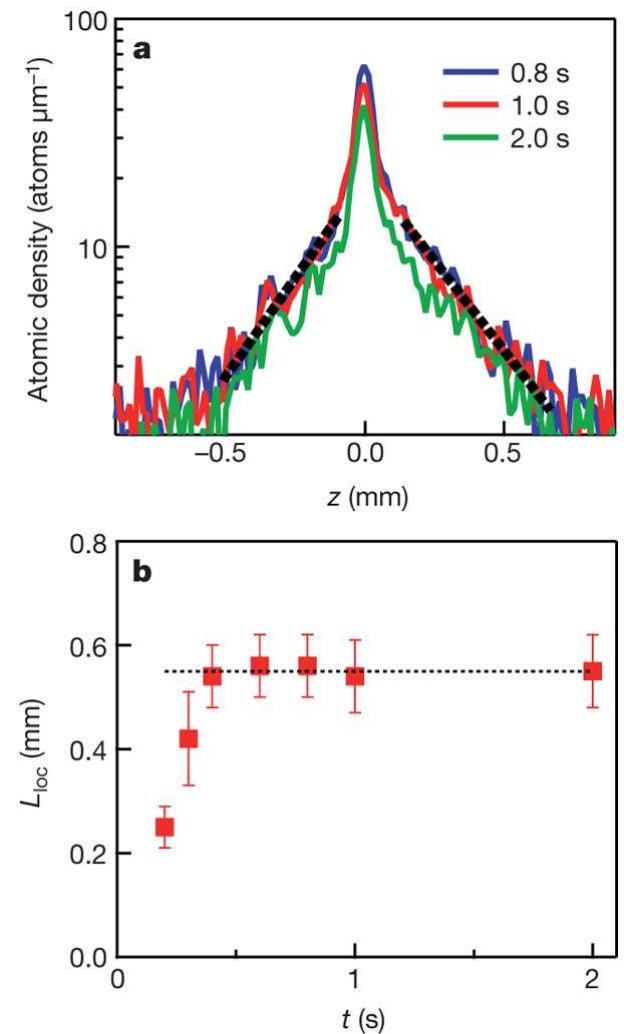
GaAs wires

Gershenson et al, PRL 97

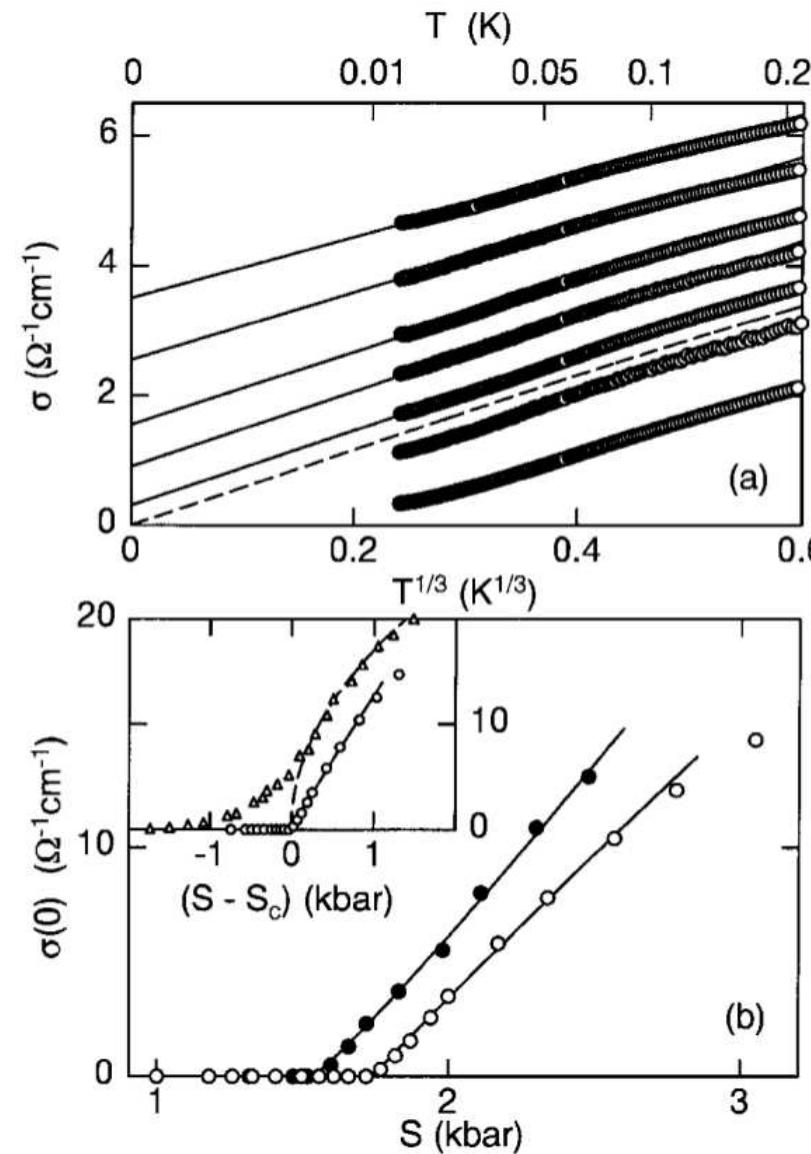
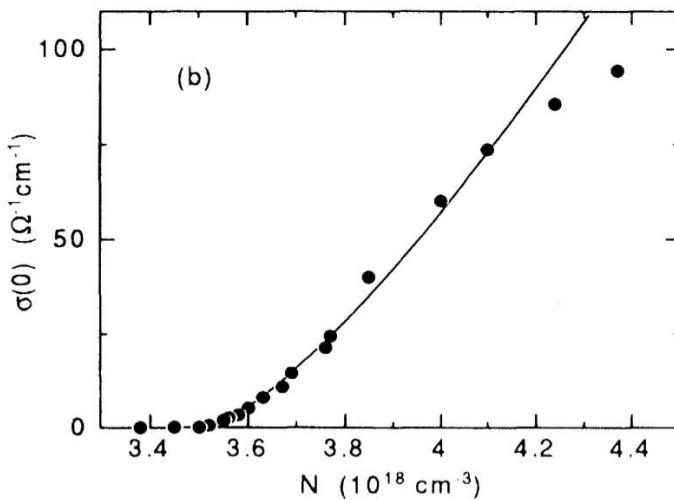
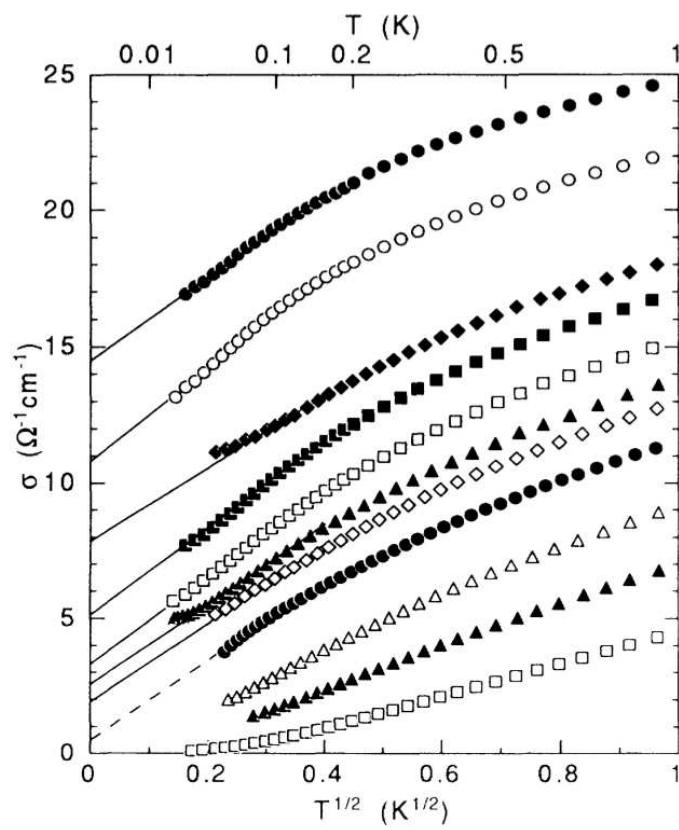
# Anderson localization of atomic Bose-Einstein condensate in 1D



Billy et al (Aspect group), Nature 2008



# 3D Anderson localization transition in Si:P



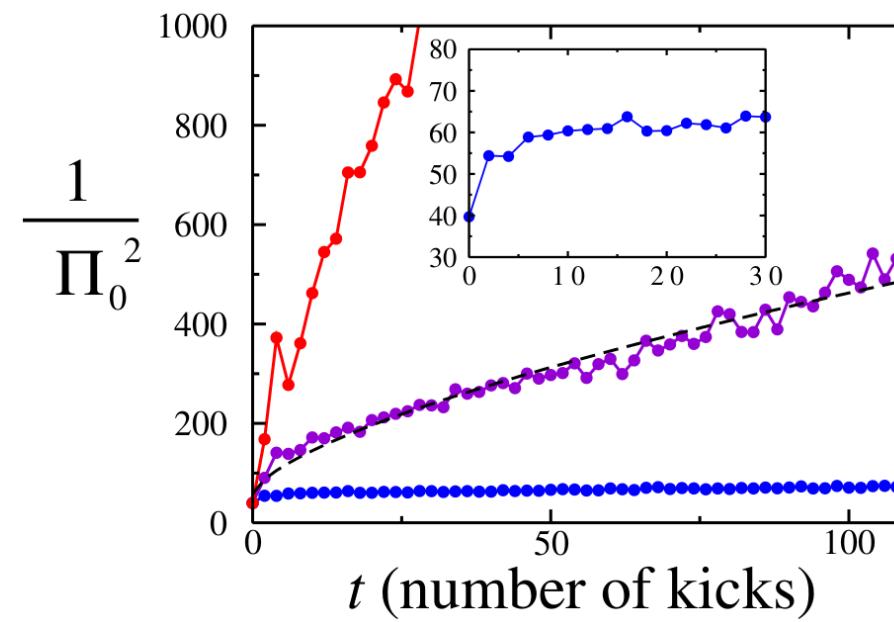
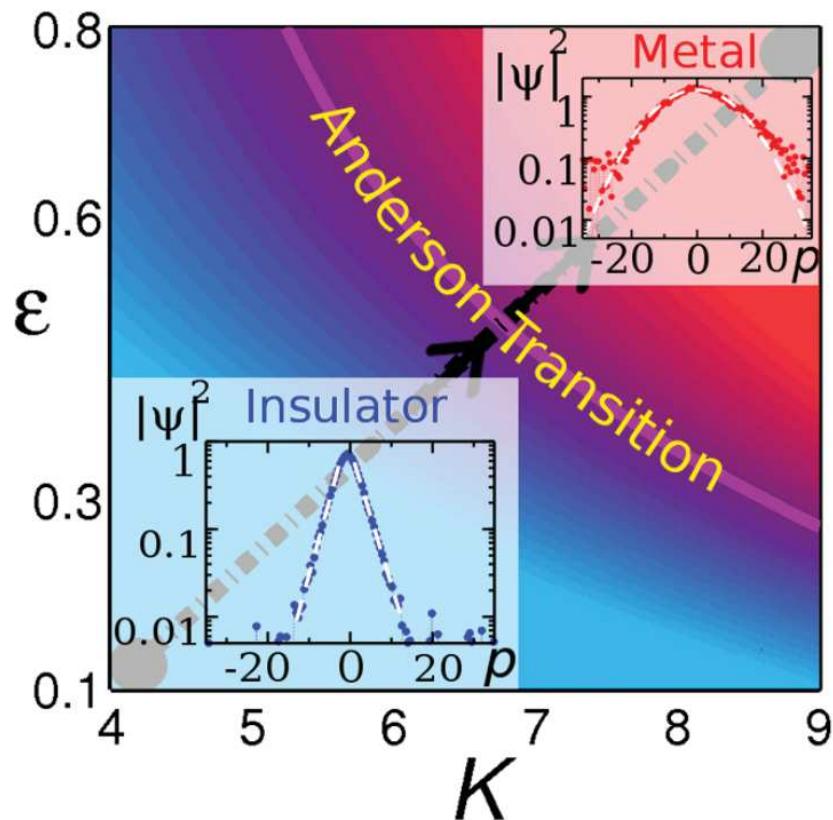
Stupp et al, PRL'93; Wafenschmidt et al, PRL'97  
(von Löhneysen group)

# 3D Anderson localization in atomic “kicked rotor”

kicked rotor       $H = \frac{p^2}{2} + K \cos x [1 + \epsilon \cos \omega_2 t \cos \omega_3 t] \sum_n \delta(t - 2\pi n/\omega_1)$

Anderson localization in momentum space. Three frequencies mimic 3D !

Experimental realization: cesium atoms exposed to a pulsed laser beam.



Chabé et al, PRL'08

# Multifractality at the Anderson transition

$$P_q = \int d^d r |\psi(r)|^{2q} \quad \text{inverse participation ratio}$$

$$\langle P_q \rangle \sim \begin{cases} L^0 & \text{insulator} \\ L^{-\tau_q} & \text{critical} \\ L^{-d(q-1)} & \text{metal} \end{cases}$$

$$\tau_q = \begin{matrix} d(q-1) & + & \Delta_q \end{matrix} \equiv D_q(q-1) \quad \text{multifractality}$$

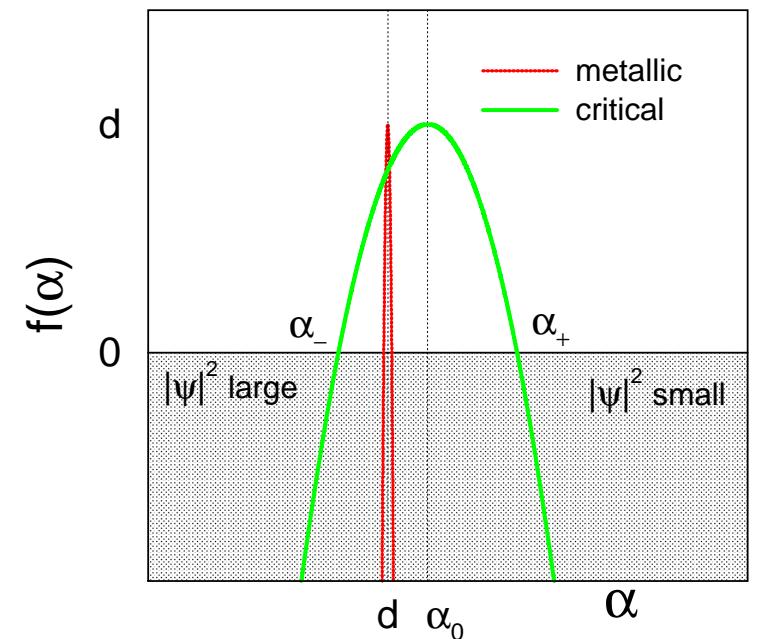
normal      anomalous

$\tau_q$  → Legendre transformation  
 → singularity spectrum  $f(\alpha)$

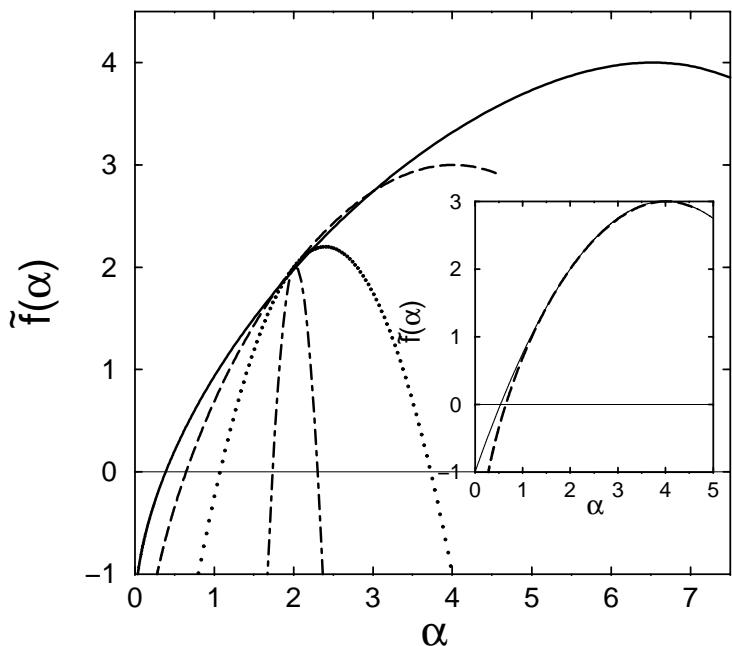
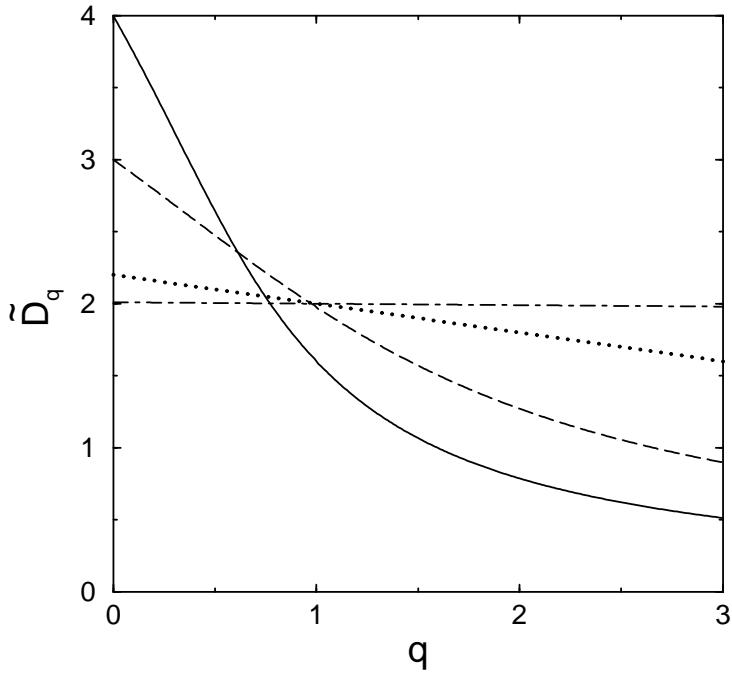
wave function statistics:

$$\mathcal{P}(\ln |\psi|^2) \sim L^{-d+f(\ln |\psi|^2 / \ln L)}$$

$L^f(\alpha)$  – measure of the set of points where  $|\psi|^2 \sim L^{-\alpha}$



# Dimensionality dependence of multifractality



Analytics (2 +  $\epsilon$ , one-loop) and numerics

$$\tau_q = (q - 1)d - q(q - 1)\epsilon + O(\epsilon^4)$$

$$f(\alpha) = d - (d + \epsilon - \alpha)^2/4\epsilon + O(\epsilon^4)$$

$d = 4$  (full)

$d = 3$  (dashed)

$d = 2 + \epsilon, \epsilon = 0.2$  (dotted)

$d = 2 + \epsilon, \epsilon = 0.01$  (dot-dashed)

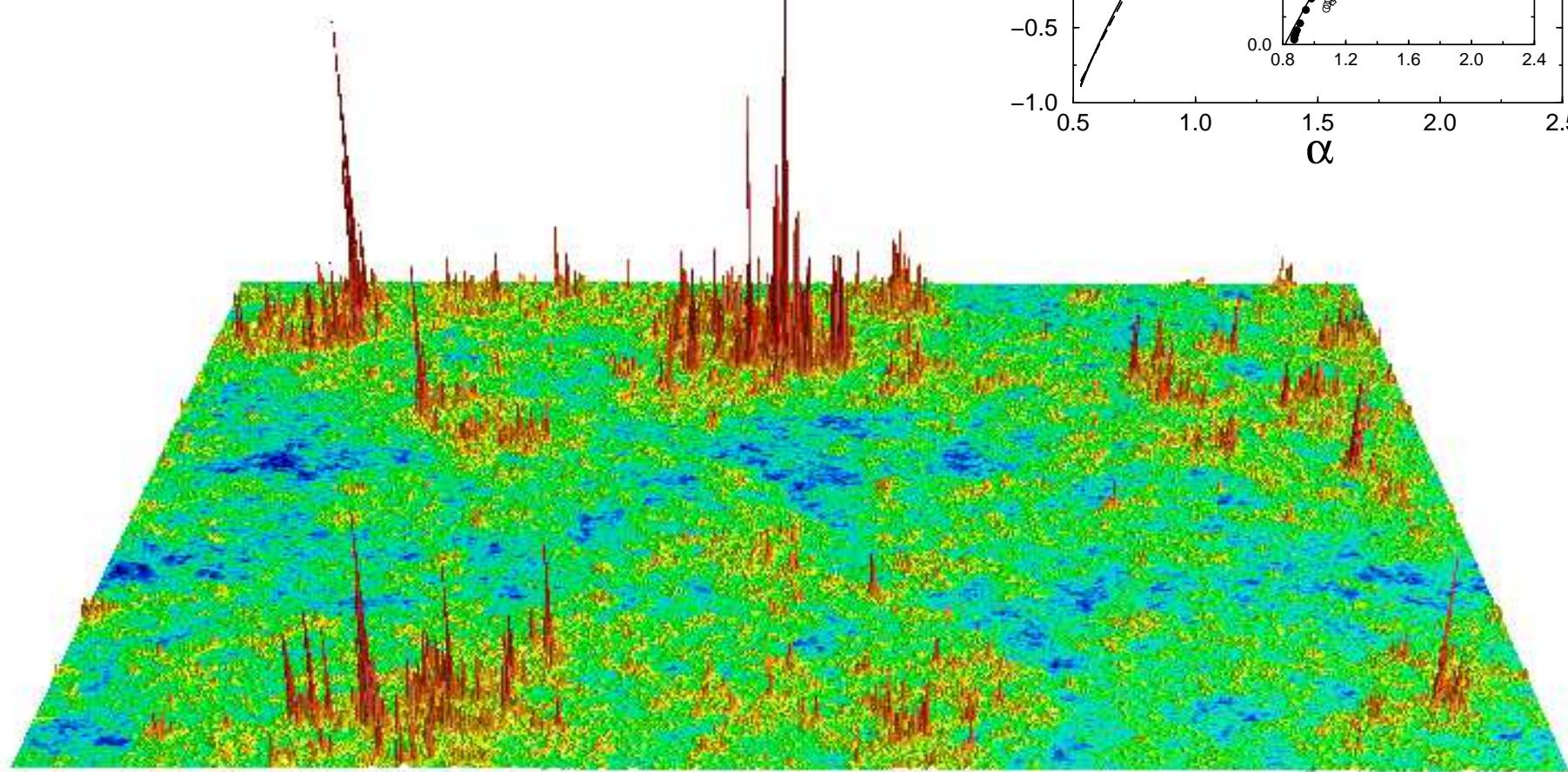
Inset:  $d = 3$  (dashed)

vs.  $d = 2 + \epsilon, \epsilon = 1$  (full)

Mildenberger, Evers, ADM '02

# Multifractality at the Quantum Hall transition

Evers, Mildenberger, ADM '01



# Symmetry of multifractal spectra

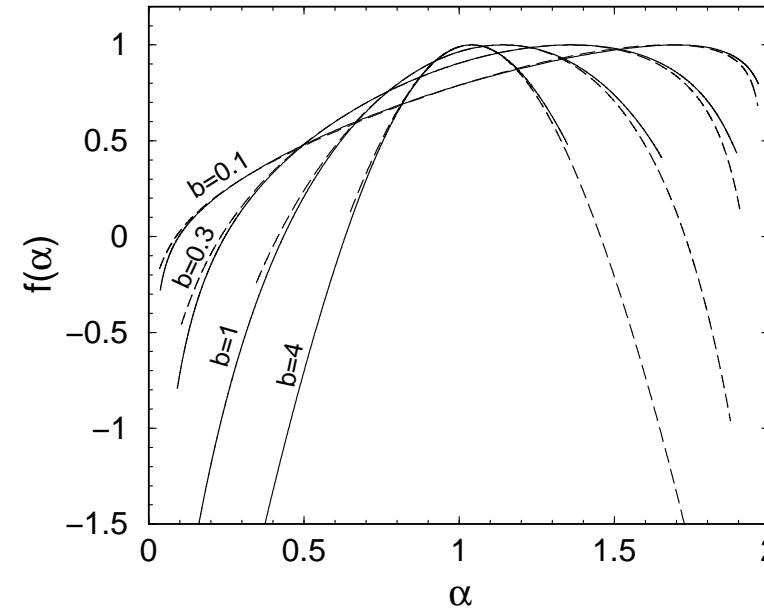
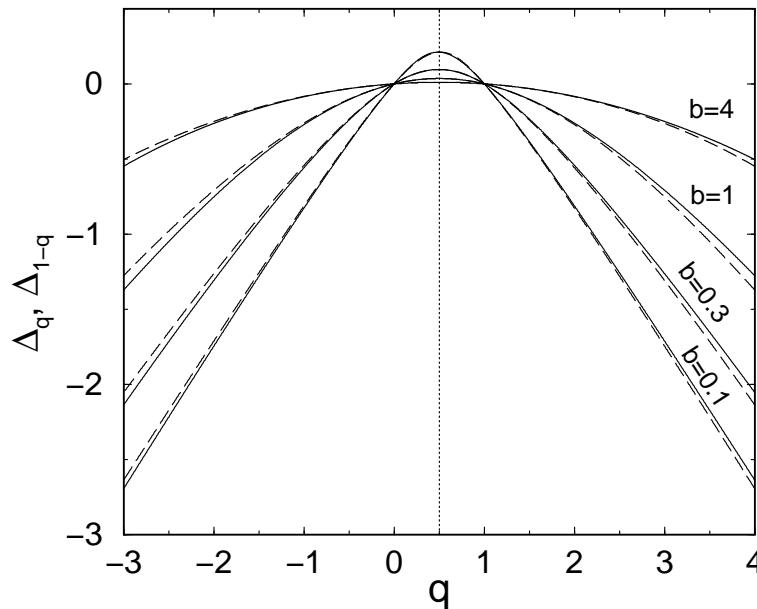
ADM, Fyodorov, Mildenberger, Evers '06

LDOS distribution in  $\sigma$ -model + universality

→ exact symmetry of the multifractal spectrum:

$$\Delta_q = \Delta_{1-q}$$

$$f(2d - \alpha) = f(\alpha) + d - \alpha$$



→ probabilities of unusually large  
and unusually small  $|\psi^2(r)|$  are related !

## Multifractality: Generalizations

- Symmetry of multifractal spectra as a consequence of invariance of the  $\sigma$  model correlation functions with respect to **Weyl group** of the  $\sigma$  model target space;

generalization to **unconventional symmetry classes**

Gruzberg, Ludwig, ADM, Zirnbauer PRL'11

- generalization on **full set of composite operators**,  
i.e. also on subleading ones.

Gruzberg, ADM, Zirnbauer, PRB'13

Important example:

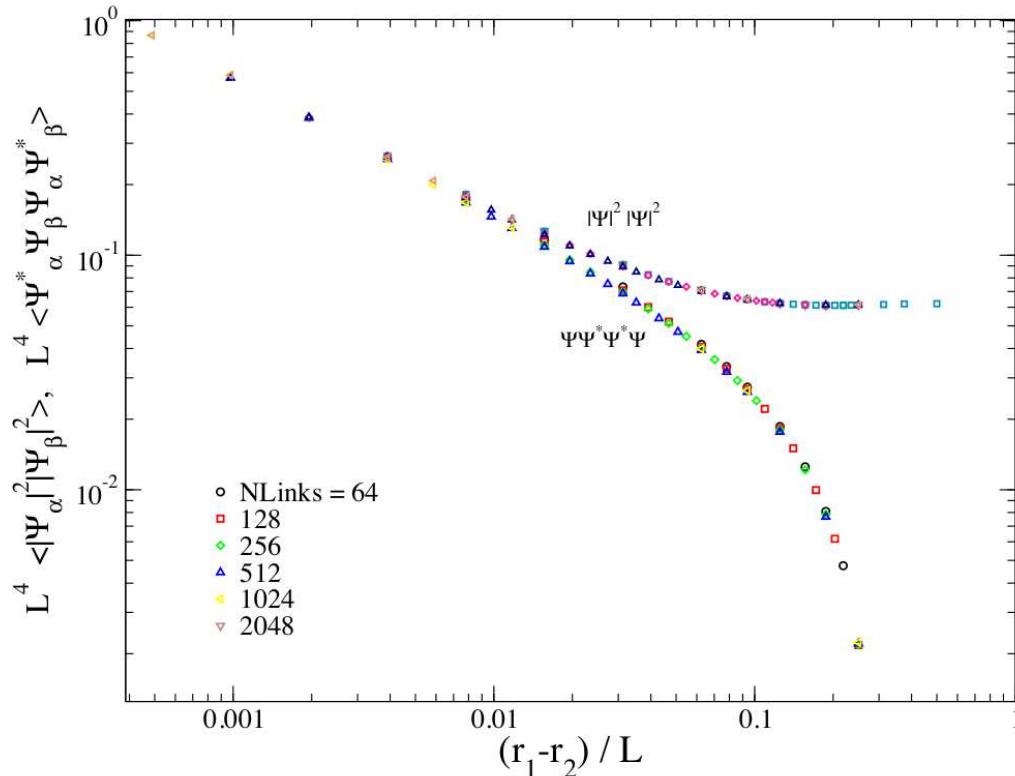
$$A_2 = V^2 |\psi_1(r_1)\psi_2(r_2) - \psi_1(r_2)\psi_2(r_1)|^2$$

↔ **Hartree-Fock matrix element of e-e interaction**

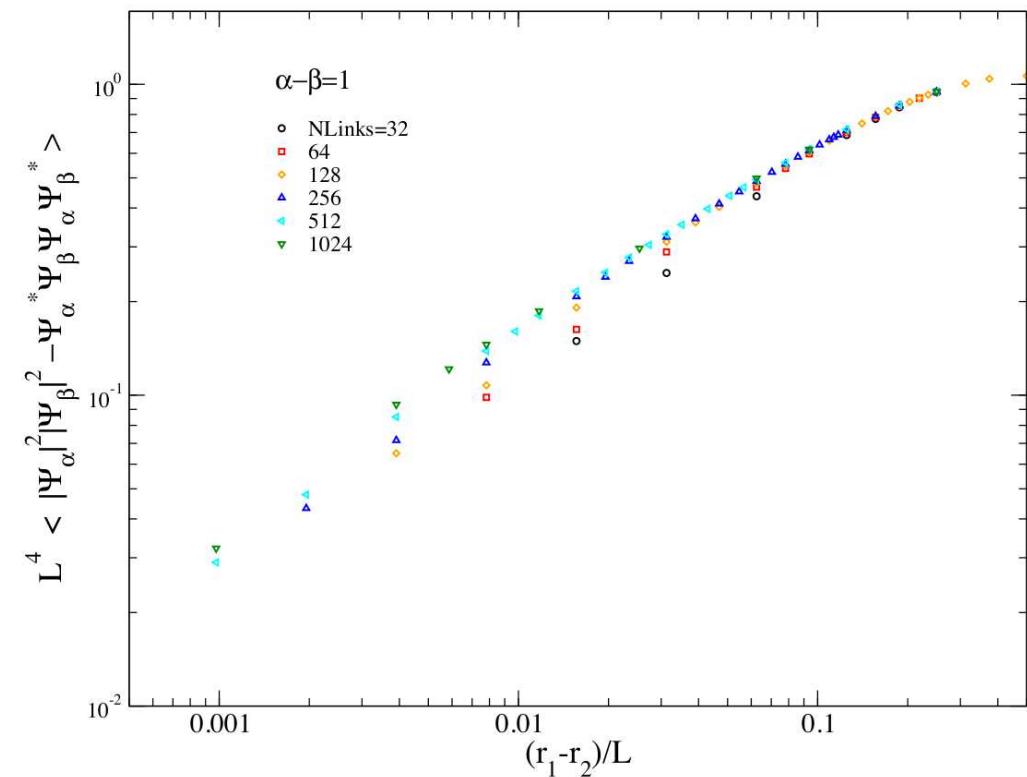
scaling:  $\langle A_2^q \rangle \propto L^{-\Delta_q^{(2)}}$

symmetry:  $\Delta_q^{(2)} = \Delta_{2-q}^{(2)}$

# Interaction scaling at criticality



Hartree, Fock  
enhanced by multifractality  
exponent  $\Delta_2 \simeq -0.52 < 0$



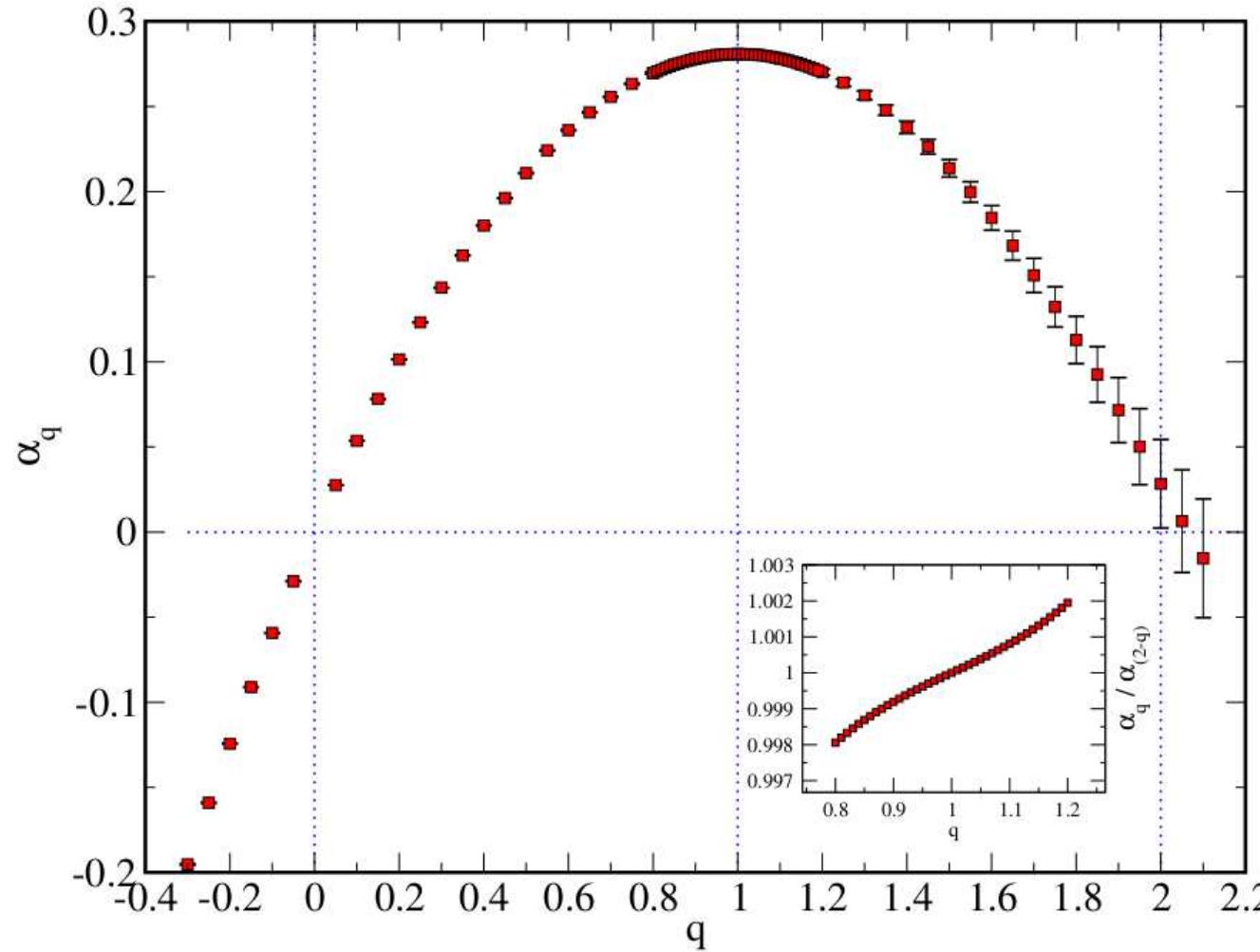
Hartree – Fock  
suppressed by multifractality  
exponent  $\Delta_1^{(2)} \simeq 0.62 > 0$

Burmistrov, Bera, Evers, Gornyi, ADM, Annals Phys. 326, 1457 (2011)

→ Temperature scaling at quantum Hall and metal-insulator transitions with short-range interaction

# Multifractal spectrum of $A_2$ at quantum Hall transition

Numerical data: Bera, Evers, unpublished

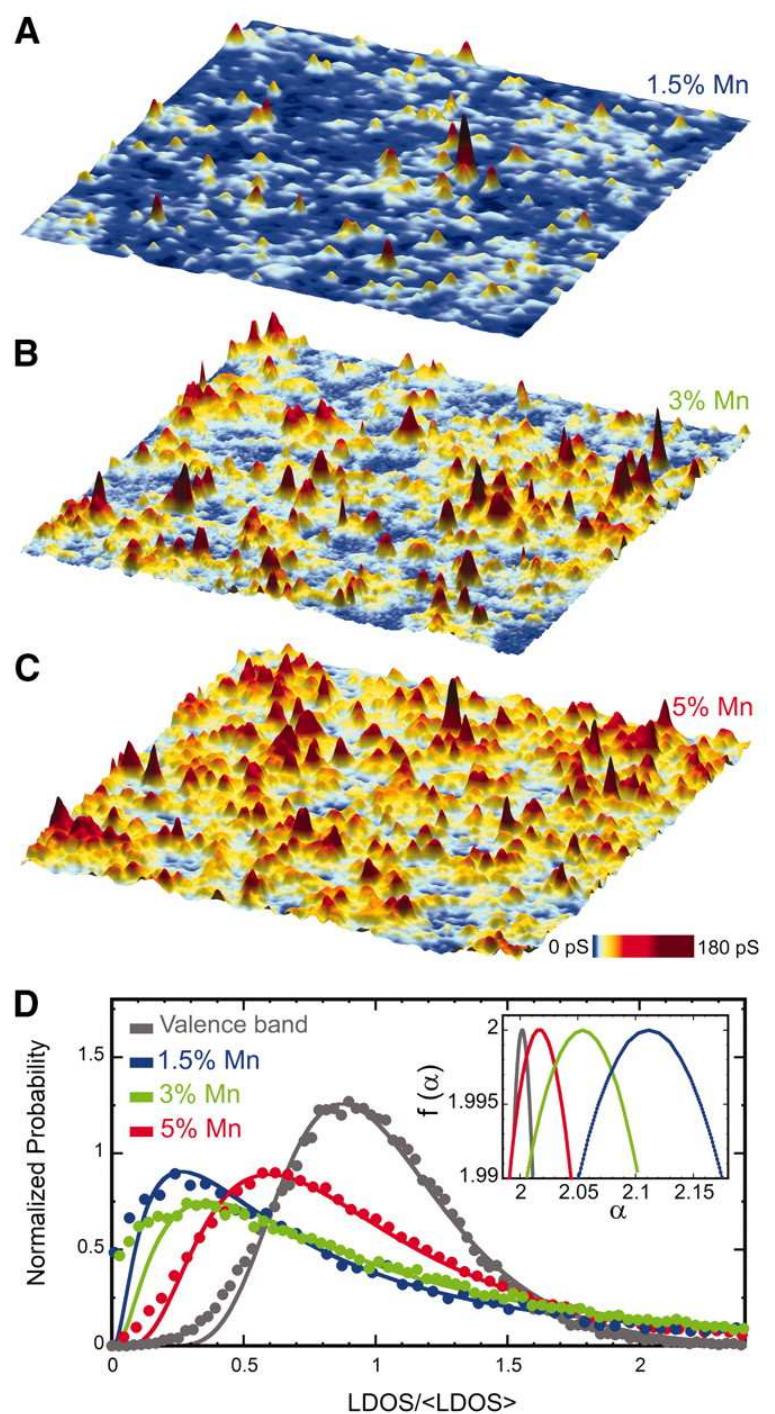


Confirms the symmetry  $q \longleftrightarrow 2 - q$

# Multifractality: Experiment I

Local DOS fluctuations  
near metal-insulator transition  
in  $\text{Ga}_{1-x}\text{Mn}_x\text{As}$

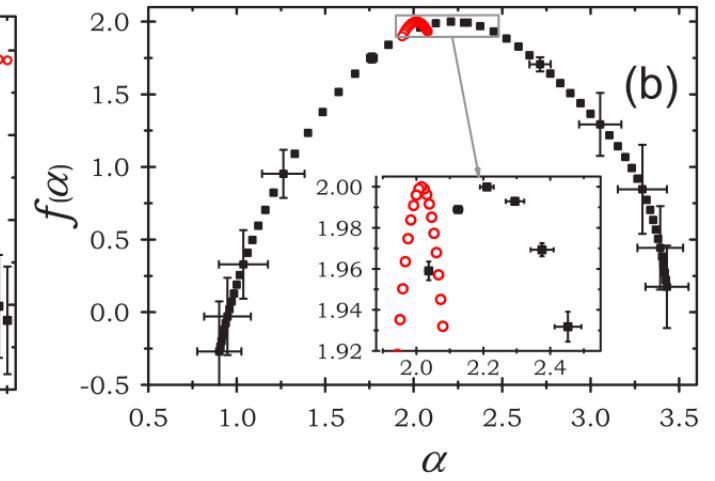
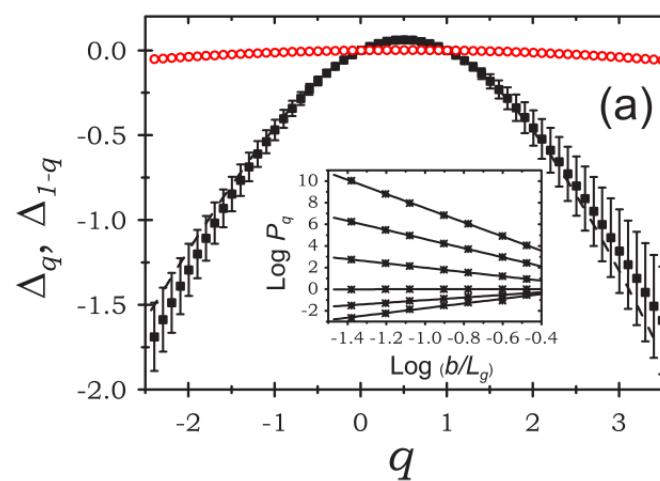
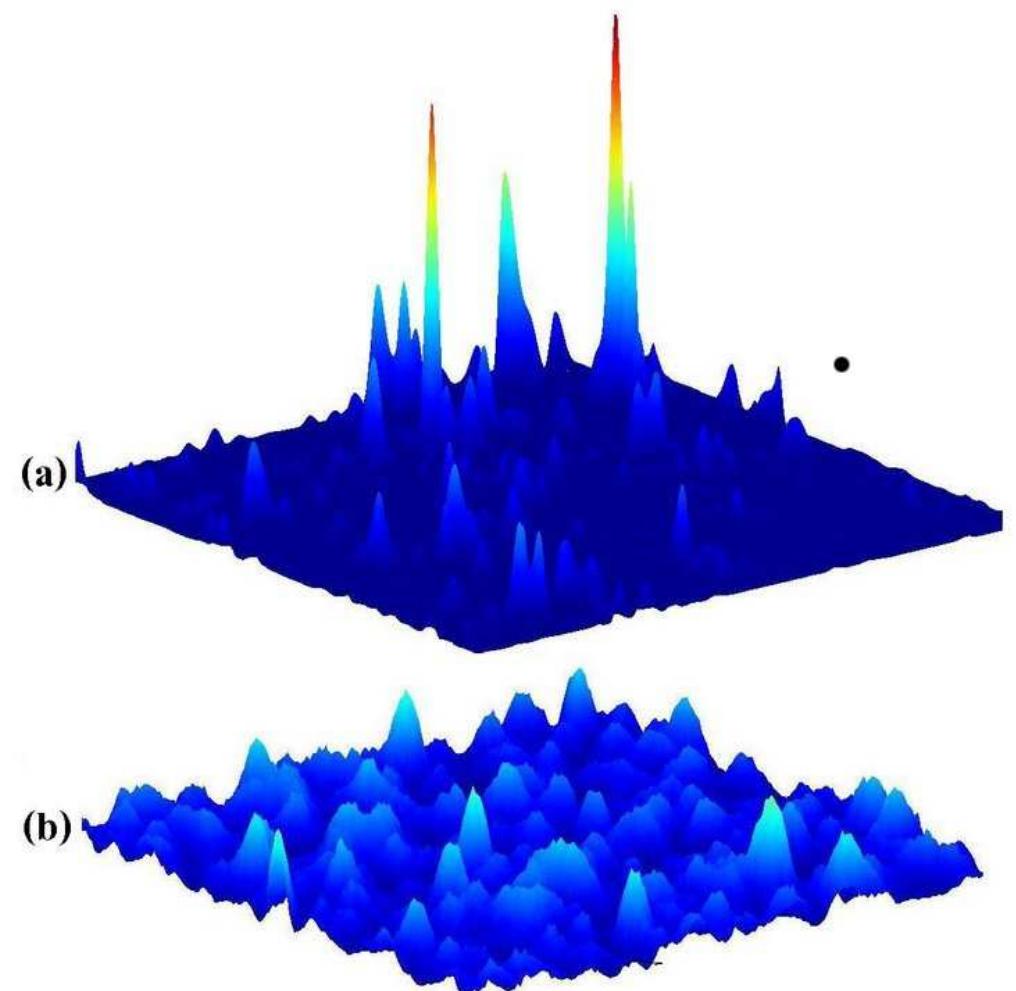
Richardella,...,Yazdani, Science '10



# Multifractality: Experiment II

Ultrasound speckle in a system  
of randomly packed Al beads

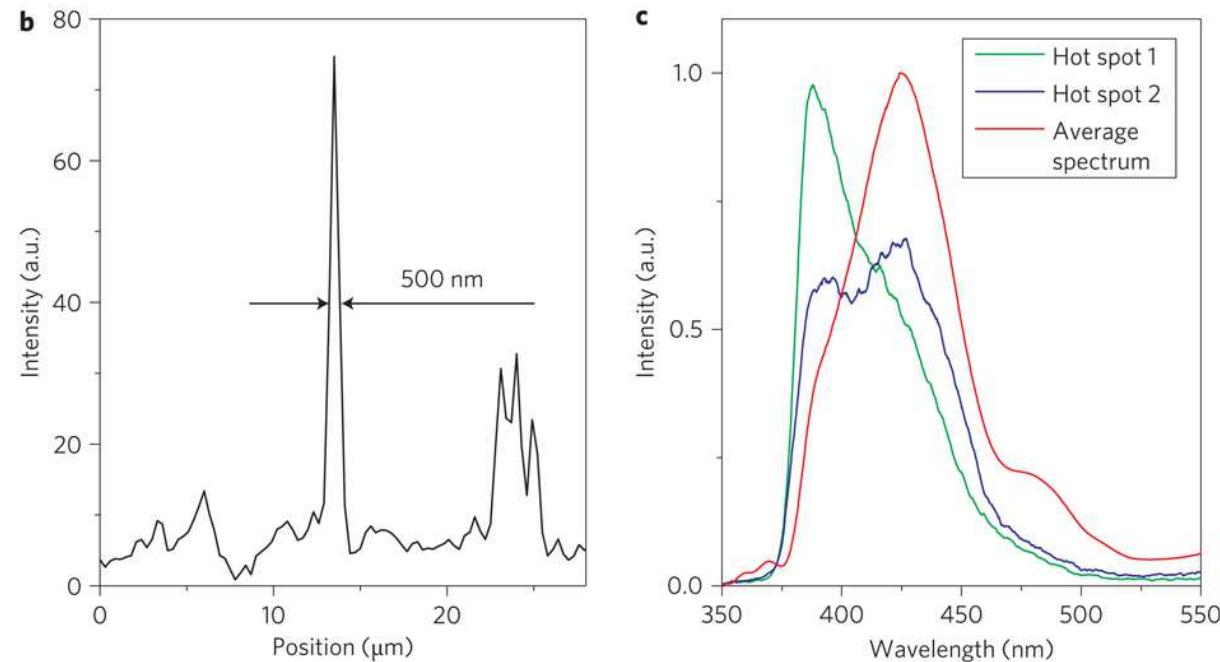
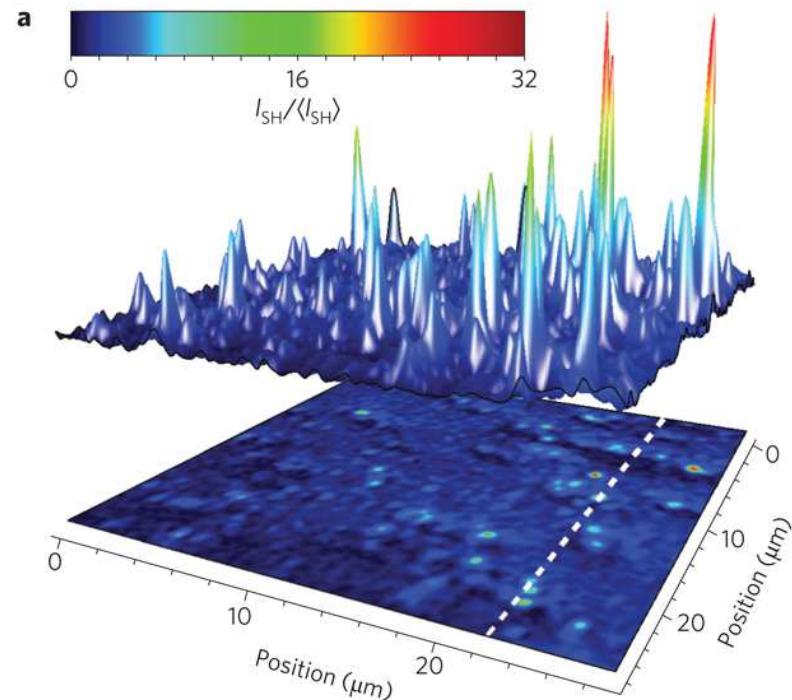
Faez, Strybulevich, Page,  
Lagendijk, van Tiggelen, PRL'09



# Multifractality: Experiment III

Localization of light  
in an array of dielectric  
nano-needles

Mascheck et al,  
Nature Photonics '12



# Disordered electronic systems: Symmetry classification

Altland, Zirnbauer '97

## Conventional (Wigner-Dyson) classes

	T	spin	rot.	symbol
GOE	+	+		AI
GUE	-	+/-		A
GSE	+	-		AII

## Chiral classes

	T	spin	rot.	symbol
ChOE	+	+		BDI
ChUE	-	+/-		AIII
ChSE	+	-		CII

$$H = \begin{pmatrix} 0 & t \\ t^\dagger & 0 \end{pmatrix}$$

## Bogoliubov-de Gennes classes

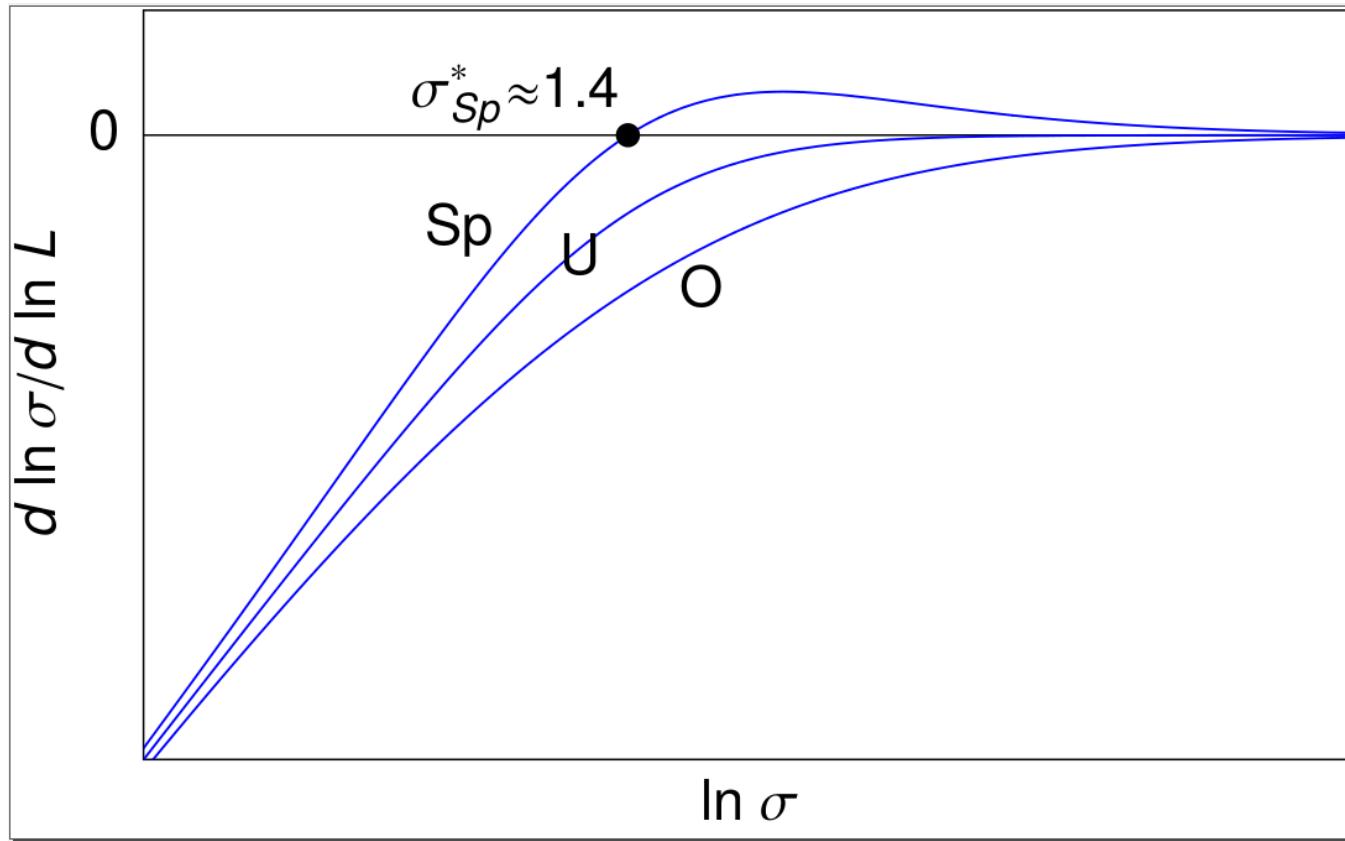
	T	spin	rot.	symbol
	+	+		CI
	-	+		C
	+	-		DIII
	-	-		D

$$H = \begin{pmatrix} h & \Delta \\ -\Delta^* & -h^T \end{pmatrix}$$

# Disordered electronic systems: Symmetry classification

Ham. class	RMT	T	S	compact symmetric space	non-compact symmetric space	$\sigma$ -model B F	$\sigma$ -model compact sector $\mathcal{M}_F$
<b>Wigner-Dyson classes</b>							
A	GUE	—	±	$U(N)$	$GL(N, \mathbb{C})/U(N)$	AIII AIII	$U(2n)/U(n) \times U(n)$
AI	GOE	+	+	$U(N)/O(N)$	$GL(N, \mathbb{R})/O(N)$	BDI CII	$Sp(4n)/Sp(2n) \times Sp(2n)$
AII	GSE	+	—	$U(2N)/Sp(2N)$	$U^*(2N)/Sp(2N)$	CII BDI	$O(2n)/O(n) \times O(n)$
<b>chiral classes</b>							
AIII	chGUE	—	±	$U(p+q)/U(p) \times U(q)$	$U(p, q)/U(p) \times U(q)$	A A	$U(n)$
BDI	chGOE	+	+	$SO(p+q)/SO(p) \times SO(q)$	$SO(p, q)/SO(p) \times SO(q)$	AI AII	$U(2n)/Sp(2n)$
CII	chGSE	+	—	$Sp(2p+2q)/Sp(2p) \times Sp(2q)$	$Sp(2p, 2q)/Sp(2p) \times Sp(2q)$	AII AI	$U(n)/O(n)$
<b>Bogoliubov - de Gennes classes</b>							
C		—	+	$Sp(2N)$	$Sp(2N, \mathbb{C})/Sp(2N)$	DIII CI	$Sp(2n)/U(n)$
CI		+	+	$Sp(2N)/U(N)$	$Sp(2N, \mathbb{R})/U(N)$	D C	$Sp(2n)$
BD		—	—	$SO(N)$	$SO(N, \mathbb{C})/SO(N)$	CI DIII	$O(2n)/U(n)$
DIII		+	—	$SO(2N)/U(N)$	$SO^*(2N)/U(N)$	C D	$O(n)$

## Role of symmetry: 2D systems of Wigner-Dyson classes



Orthogonal and Unitary: localization;  
parametrically different localization length:  $\xi_U \gg \xi_O$

Symplectic: metal-insulator transition

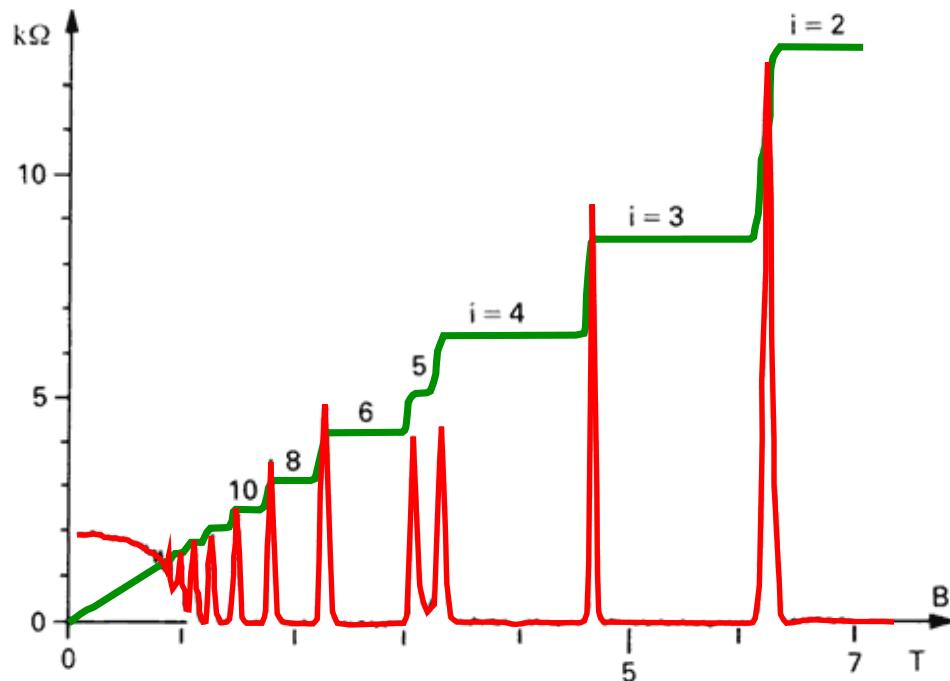
Usual realization of Sp class: spin-orbit interaction

Symmetry alone is not always sufficient to characterize the system.

There may be also a non-trivial **topology**.

It may **protect** the system from localization.

# Integer quantum Hall effect



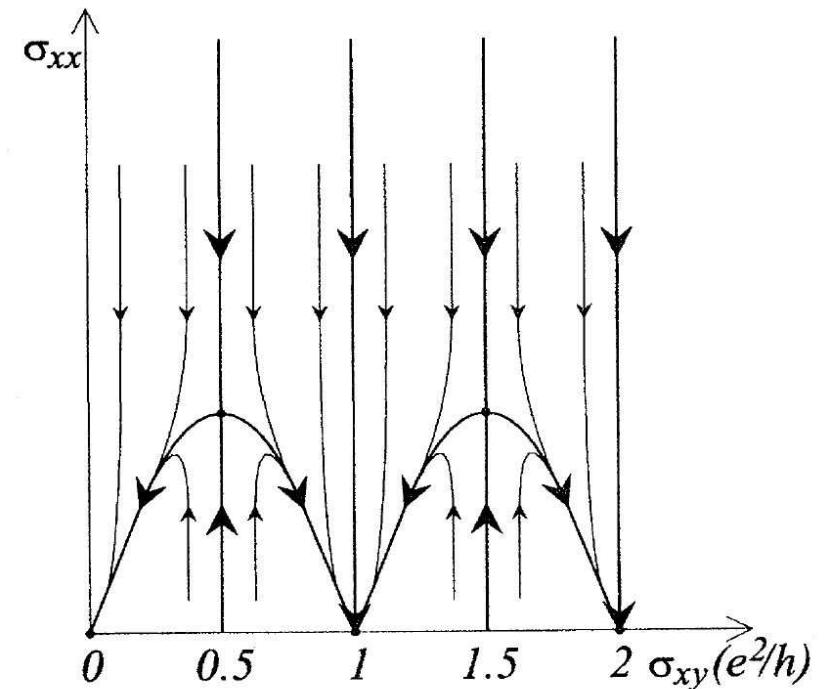
von Klitzing '80 ; Nobel Prize '85

Field theory (Pruisken):

$\sigma$ -model with topological term

$$S = \int d^2r \left\{ -\frac{\sigma_{xx}}{8} \text{Tr}(\partial_\mu Q)^2 + \frac{\sigma_{xy}}{8} \text{Tr} \epsilon_{\mu\nu} Q \partial_\mu Q \partial_\nu Q \right\}$$

QH insulators  $\rightarrow n = \dots, -2, -1, 0, 1, 2, \dots$  protected edge states  
 $\rightarrow \mathbb{Z}$  topological insulator



IQHE flow diagram

Khmelnitskii' 83, Pruisken' 84

localized

localized

critical point

# Periodic table of Topological Insulators

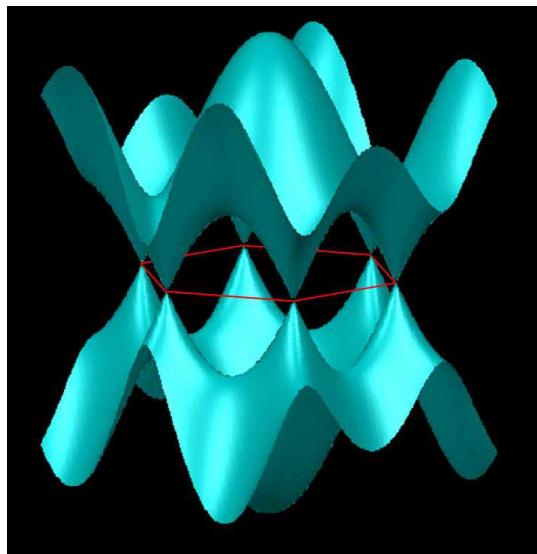
Symmetry classes					Topological insulators			
$p$	$H_p$	$R_p$	$S_p$	$\pi_0(R_p)$	d=1	d=2	d=3	d=4
0	AI	BDI	CII	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$
1	BDI	BD	AII	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0
2	BD	DIII	DIII	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0
3	DIII	AII	BD	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0
4	AII	CII	BDI	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$
5	CII	C	AI	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$
6	C	CI	CI	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$
7	CI	AI	C	0	0	0	$\mathbb{Z}$	0
$0'$	A	AIII	AIII	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$
$1'$	AIII	A	A	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0

$H_p$  – symmetry class of Hamiltonians

$R_p$  – sym. class of classifying space (of Hamiltonians with eigenvalues  $\rightarrow \pm 1$ )

$S_p$  – symmetry class of compact sector of  $\sigma$ -model manifold

## 2D massless Dirac fermions

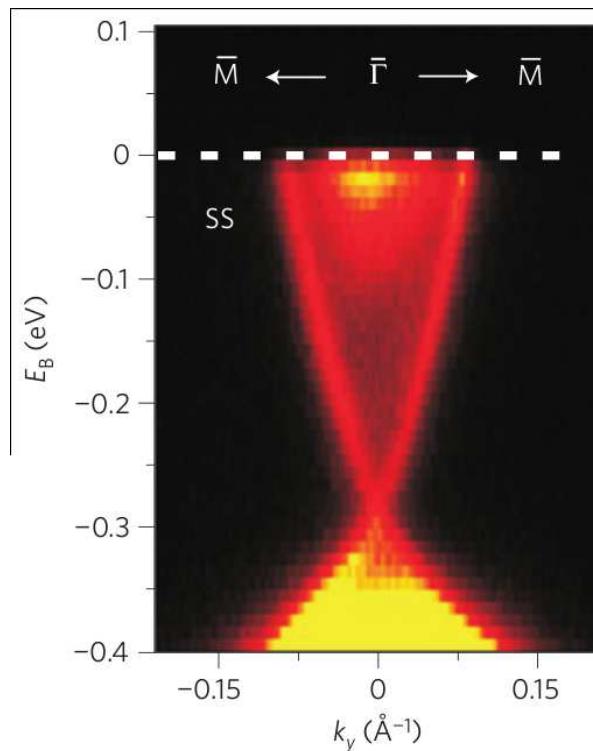


Graphene

Geim, Novoselov'04  
Nobel Prize'10

$\sigma$ -model field theory with a topological term

- Graphene: long-range disorder (no valley mixing)
- Surface states of 3D TI: no restriction on disorder range

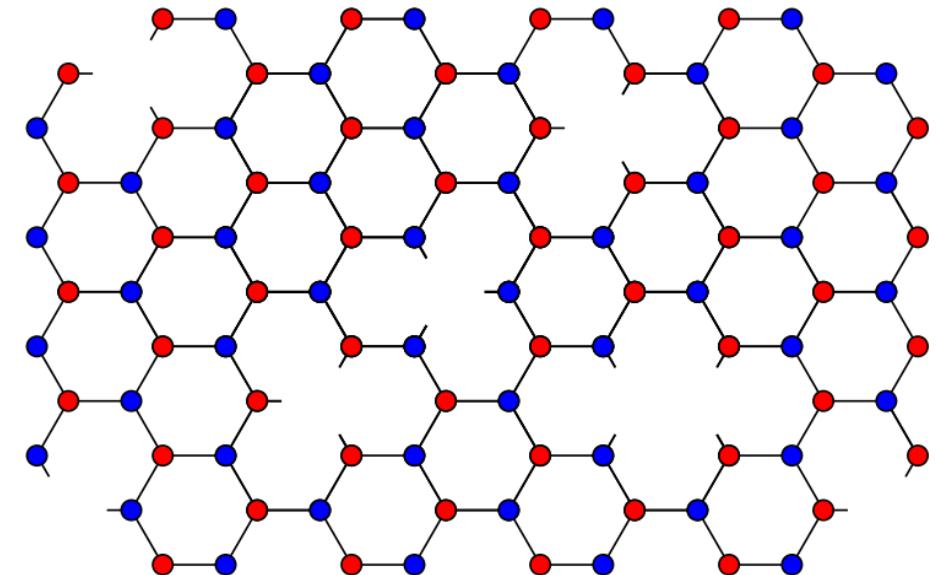
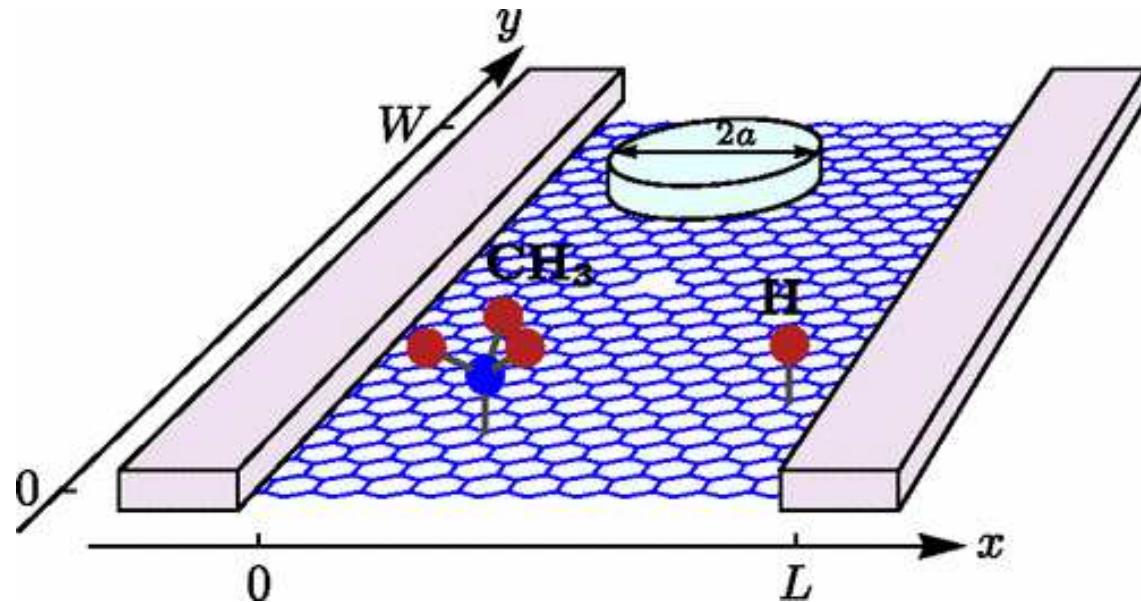


Surface of 3D topological insulators  
BiSb, BiSe, BiTe      Hasan group '08

Ostrovsky, Gornyi, ADM '07

# Role of symmetry and topology: Graphene at the Dirac point

Ostrovsky et al, PRL'10; Gattenlöchner et al, PRL'14

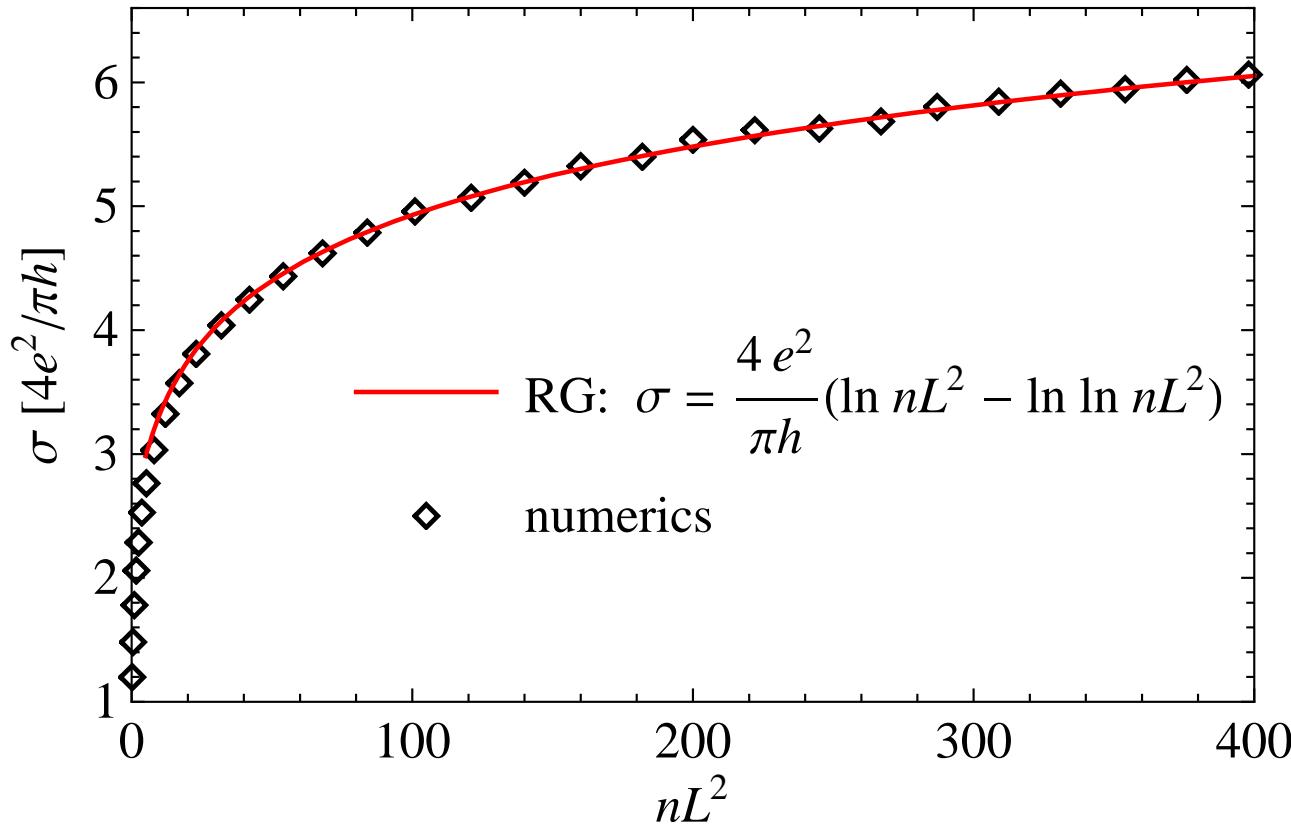


Models of scatterers:

- **scalar impurity**: smooth on atomic scale (no valley mixing)
- **resonant scalar impurity**: diverging scattering length, quasi-bound state at the Dirac point
- **adatom**: on-site potential (valley mixing)
- **vacancy**: infinitely strong on-site potential

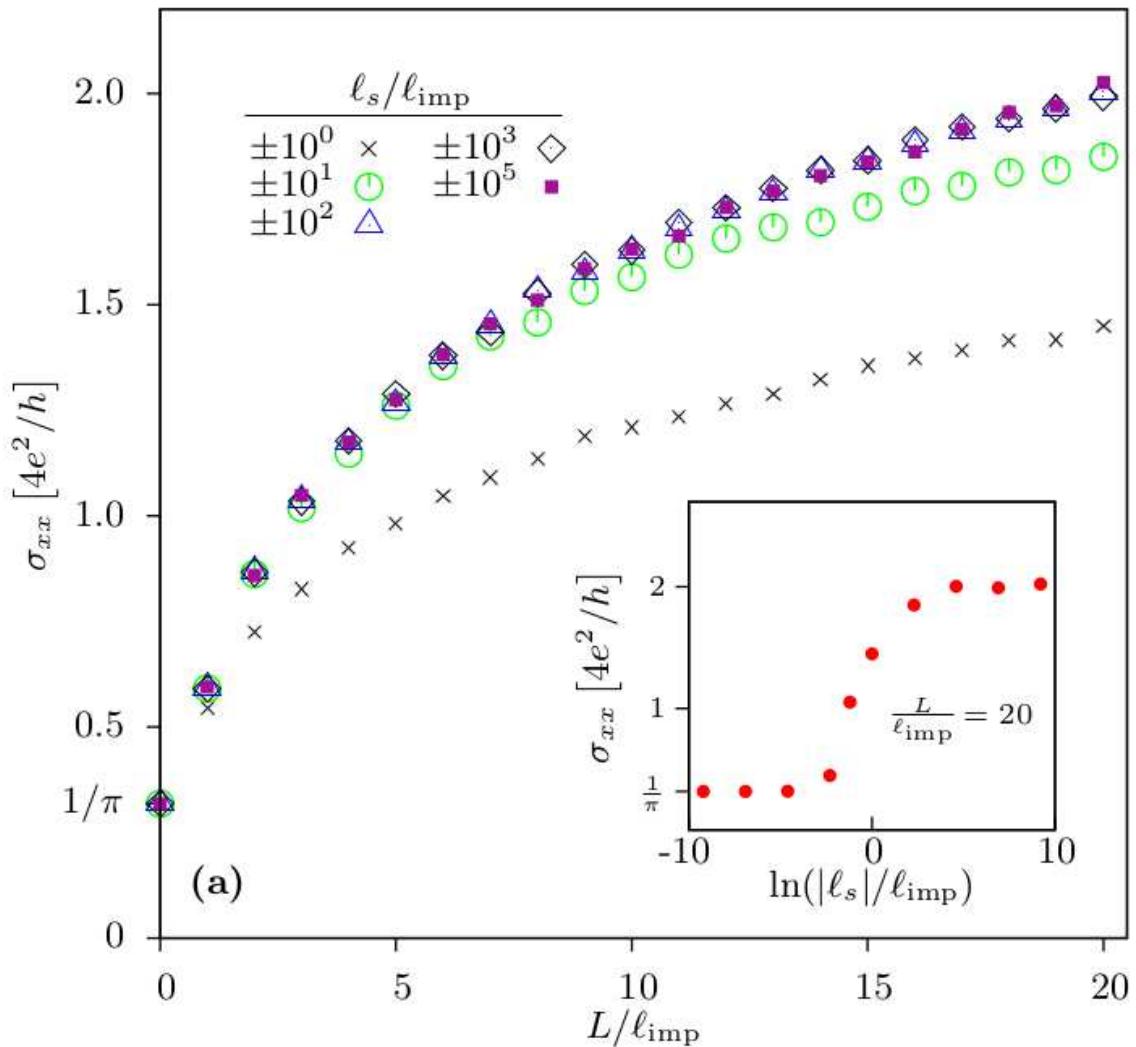
# Resonant scalar impurities ( $l_s = \infty$ )

Ostrovsky, Titov, Bera, Gornyi, ADM, PRL (2010)



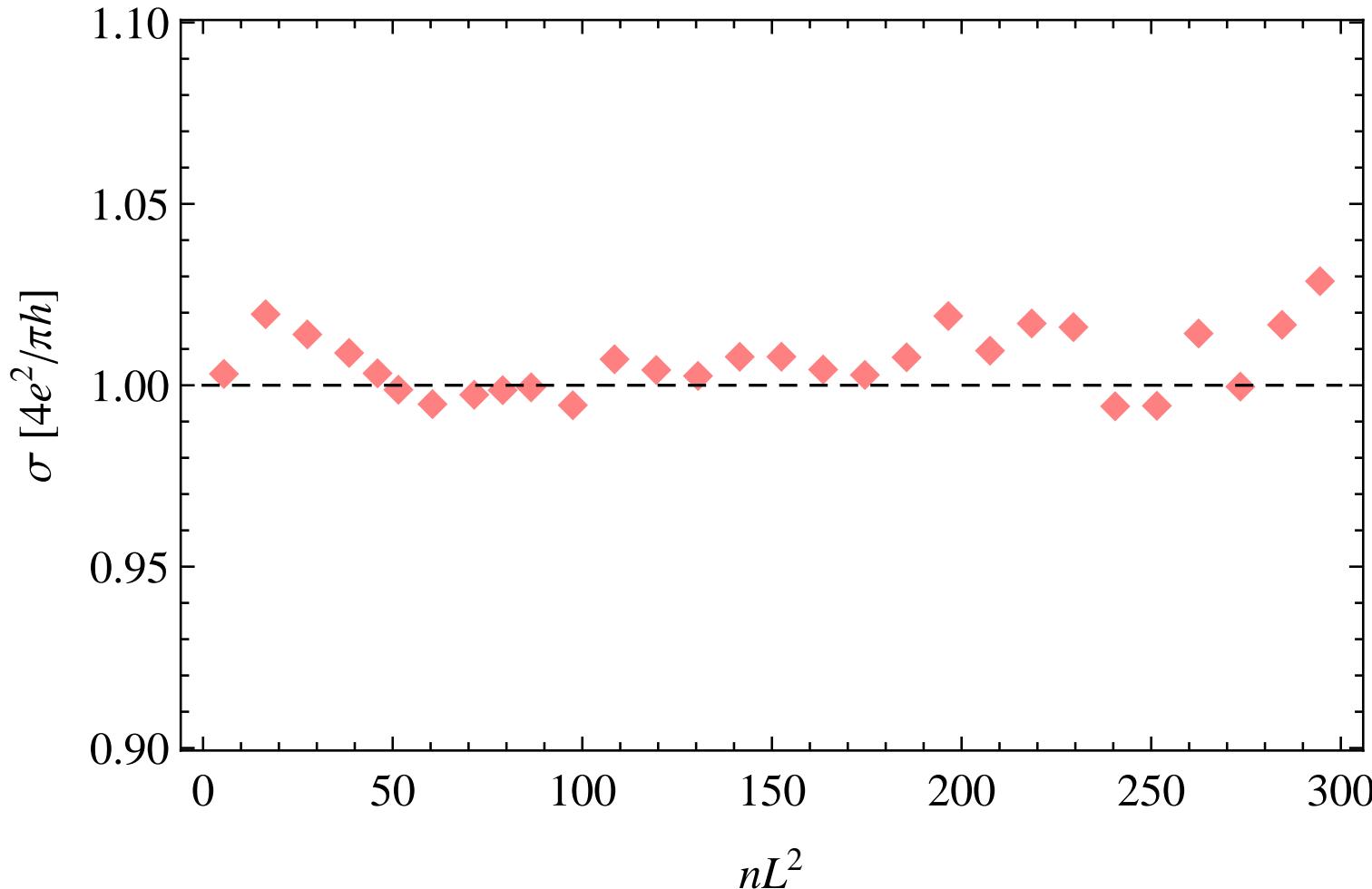
- flow towards supermetal  $\sigma \rightarrow \infty$
- agreement with  $\sigma$  model RG

# Scalar impurities (finite $l_s$ , random sign)



Large  $l_s \rightarrow$  Symmetry breaking pattern:  
DIII (with WZ term)  $\rightarrow$  AII (with  $\mathbb{Z}_2 \theta$ -term)

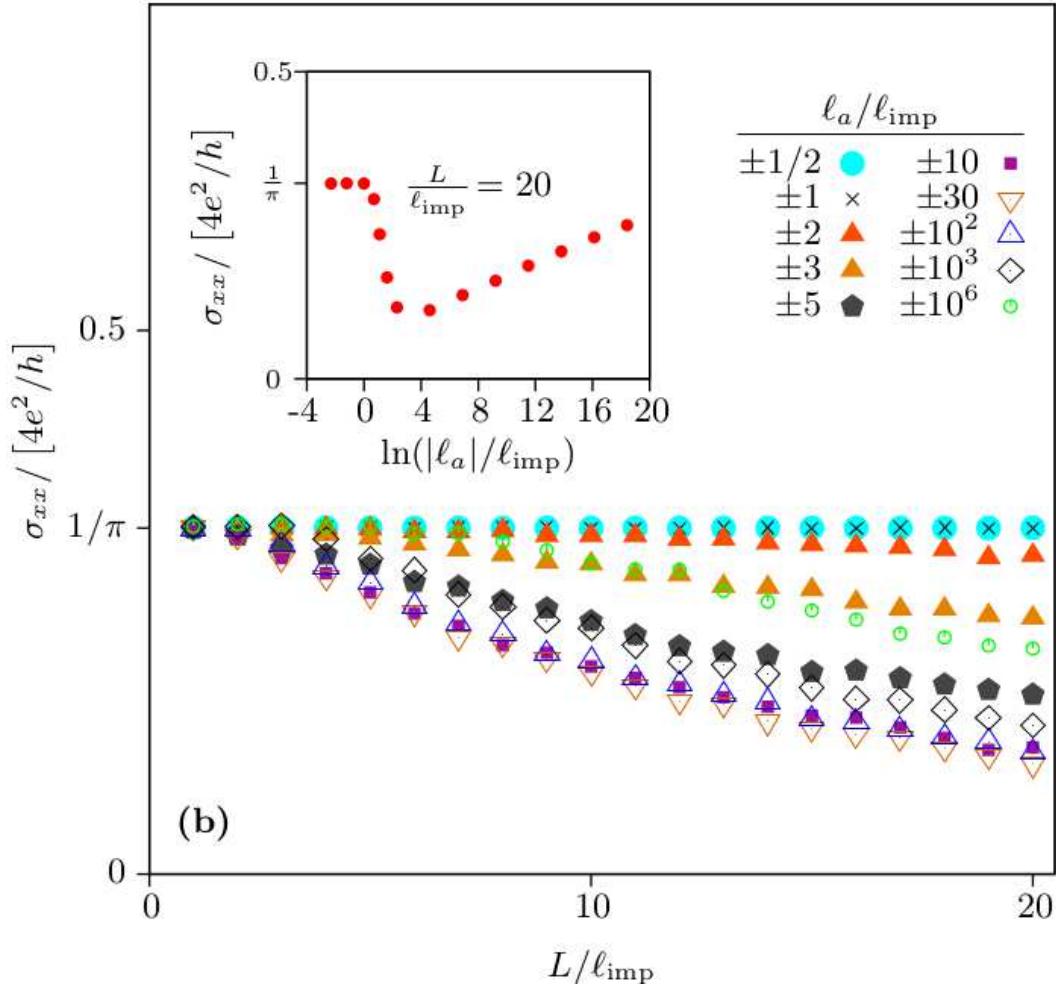
# Vacancies



symmetry class BDI (chiral orthogonal)

No localization,  $\sigma \rightarrow \text{const} \simeq \frac{4 e^2}{\pi h}$

# Adatoms (finite $l_a$ , random sign)

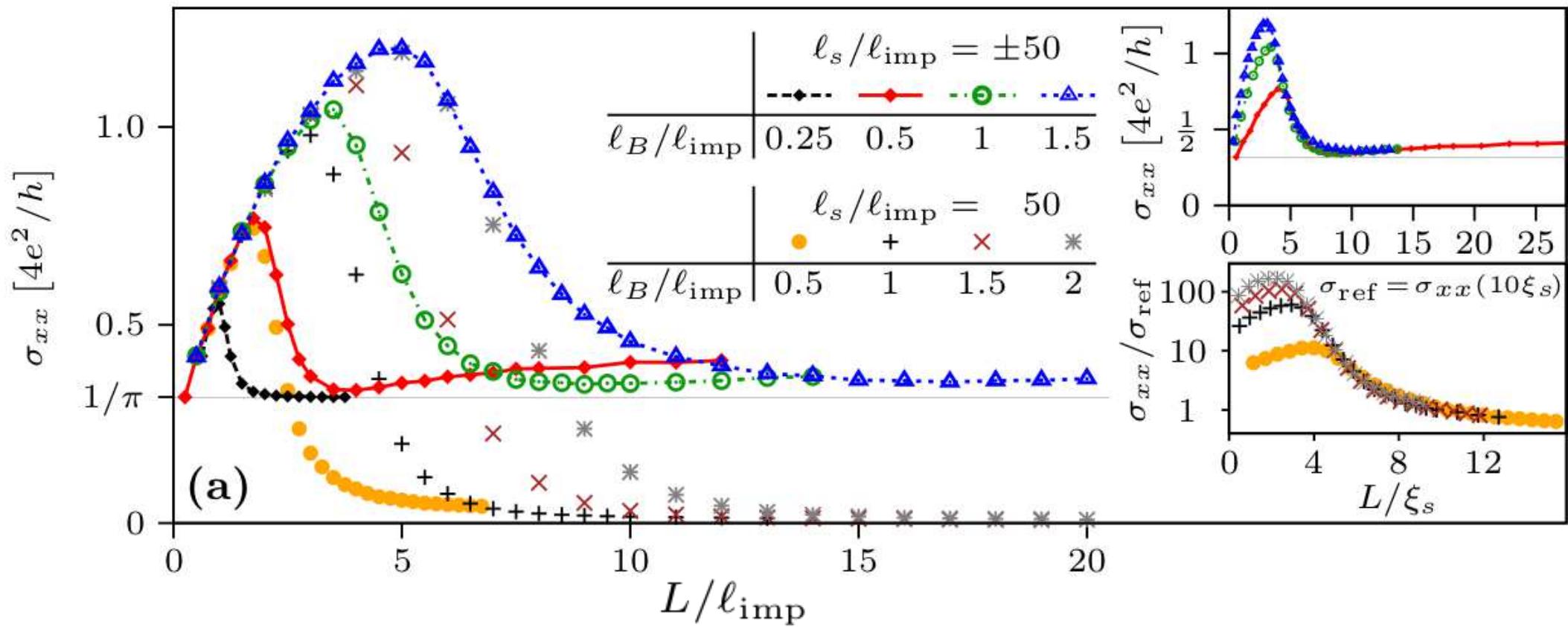


Large  $l_a \rightarrow$  Symmetry breaking pattern: BDI  $\rightarrow$  AI

Vacancies ( $l_a \rightarrow \infty$ ): finite conductivity  $\sigma \simeq \frac{4e^2}{\pi h}$  for  $L \rightarrow \infty$

Localization length  $\xi$  – non-monotonous function of  $l_a$

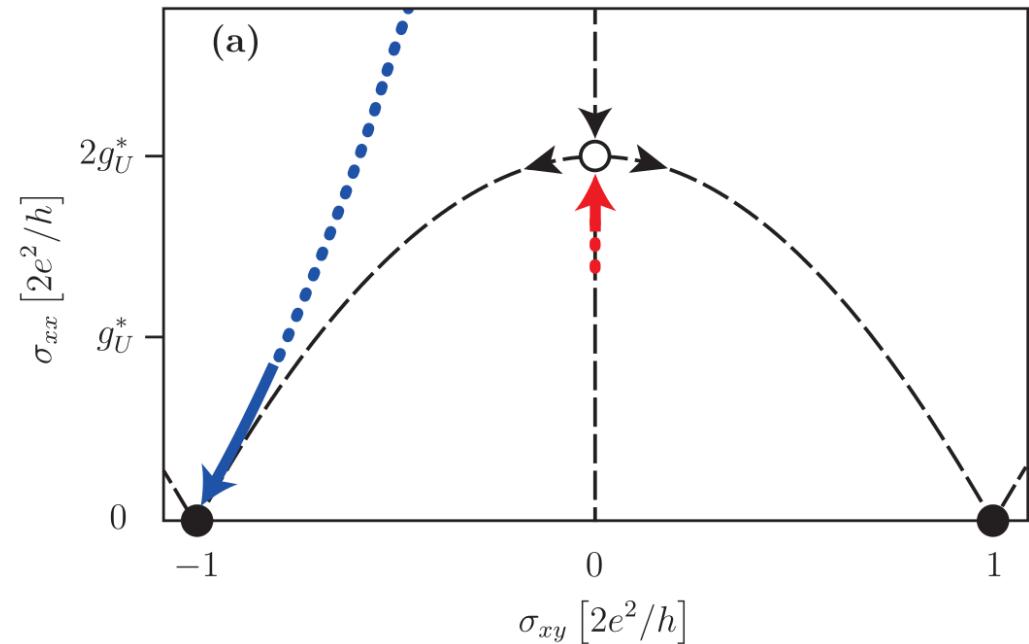
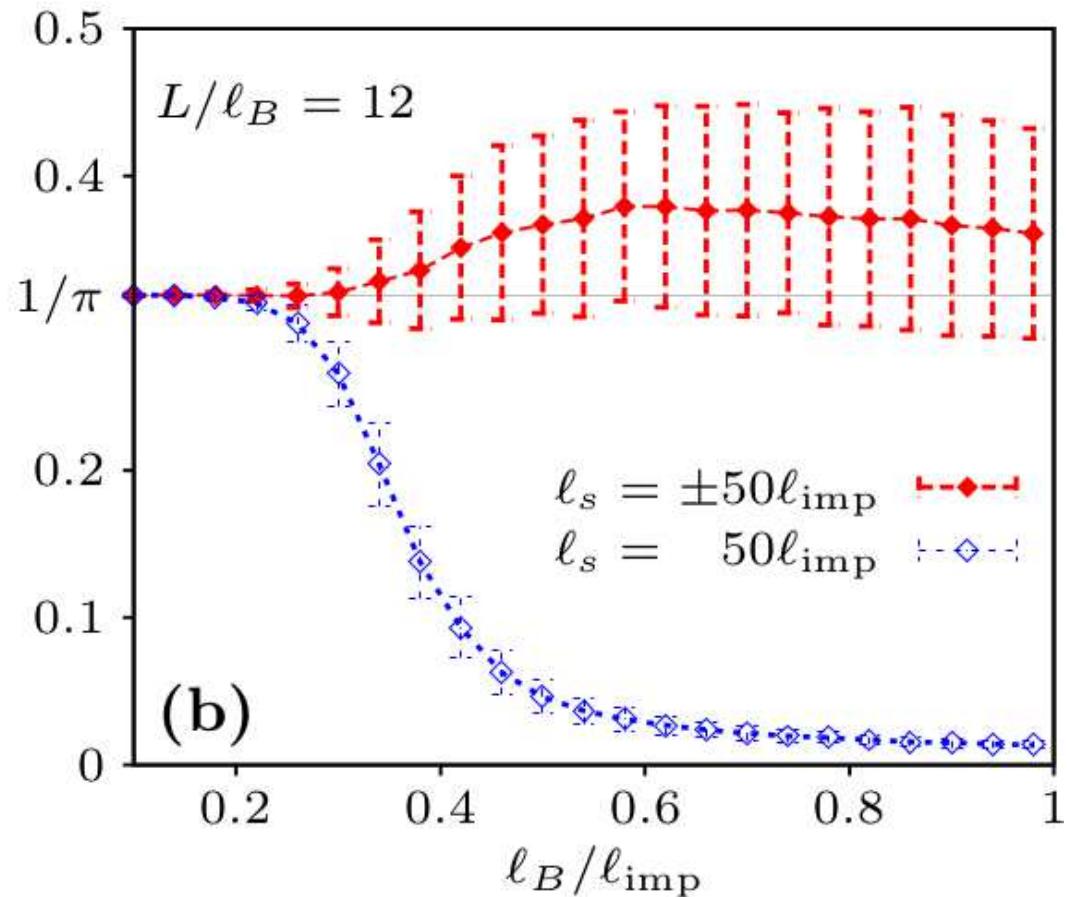
# Scalar impurities in magnetic field $B$



Symmetry breaking pattern: DIII  $\rightarrow$  AII  $\rightarrow$  A for weaker B  
and DIII  $\rightarrow$  AIII  $\rightarrow$  A for stronger B

Ultimate fixed points: Quantum Hall criticality (random sign of impurity potentials) and localization (fixed sign)

# Scalar impurities in magnetic field $B$

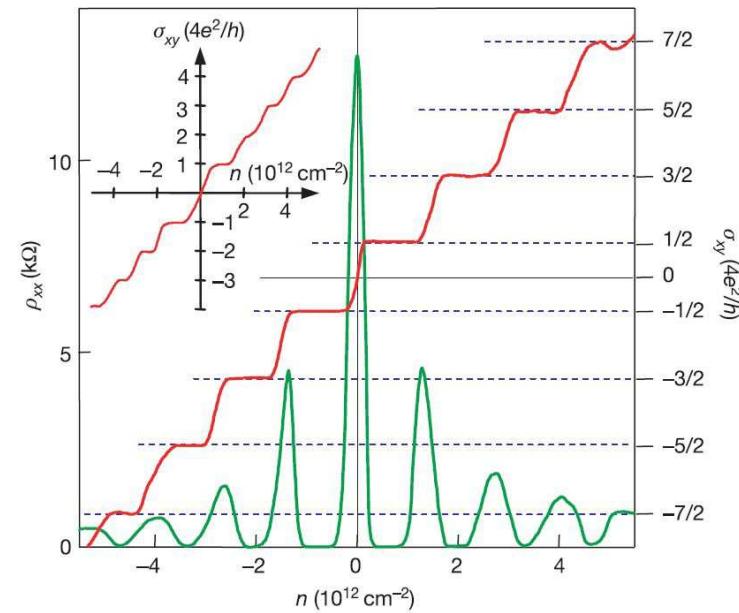
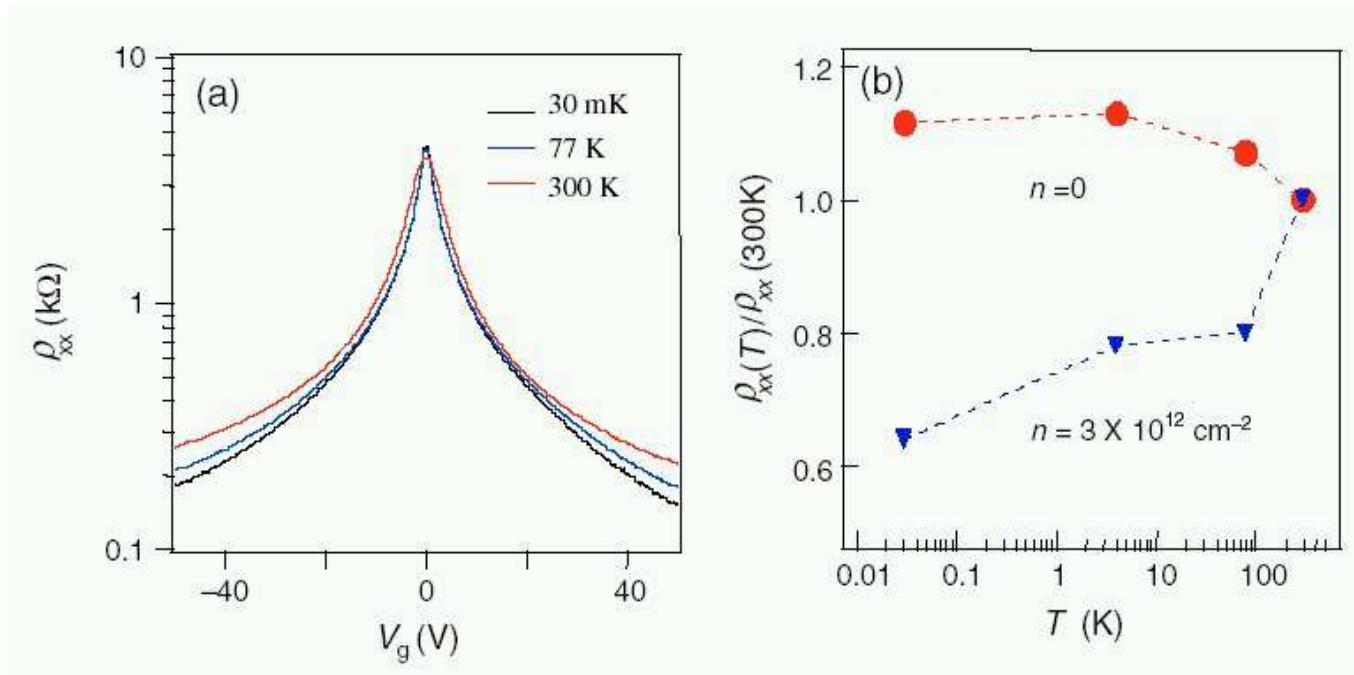


Ultimate fixed points: Quantum Hall criticality (random sign of impurity potentials) and localization (fixed sign)

Vertical bars: mesoscopic fluctuations

# Graphene: Experiments

Geim-Novoselov and Kim groups



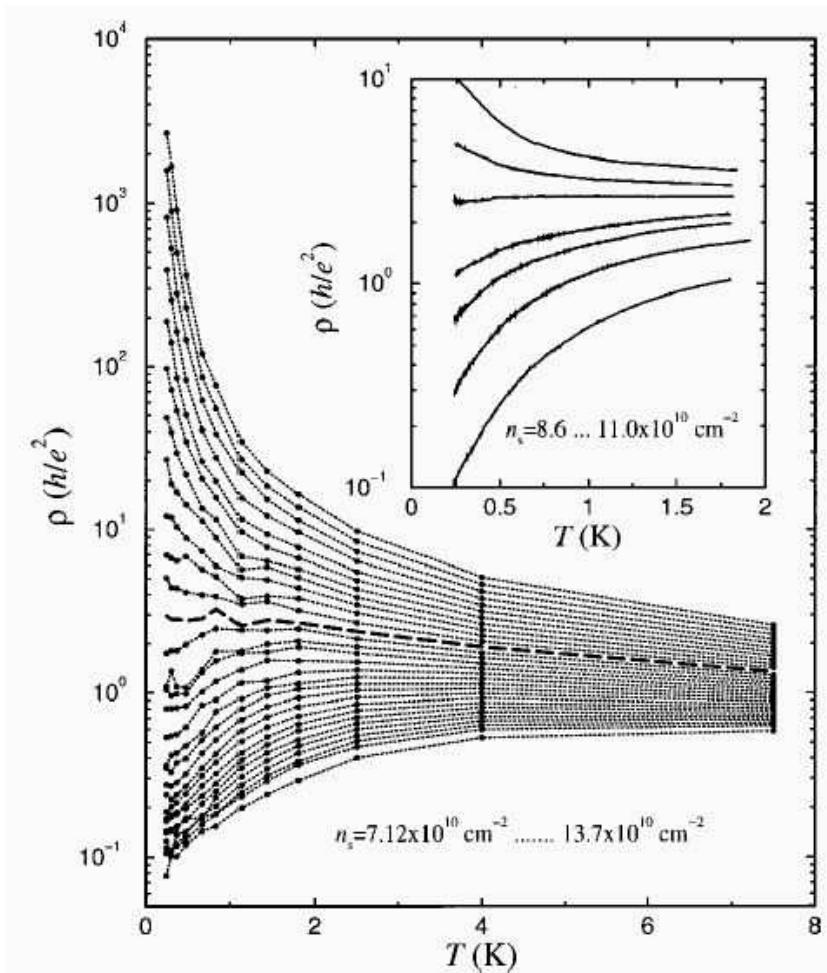
Topological terms explain unconventional properties of high-quality graphene samples:

- absence of localization at Dirac point down to very low temperatures (30 mK), although conductivity  $\simeq e^2/h$  per spin per valley
- anomalous QHE:  $\sigma_{xy} = (2n + 1) \times 2e^2/h$ ; QHE transition at  $n = 0$  (Dirac point), i.e. at  $\sigma_{xy} = 0$

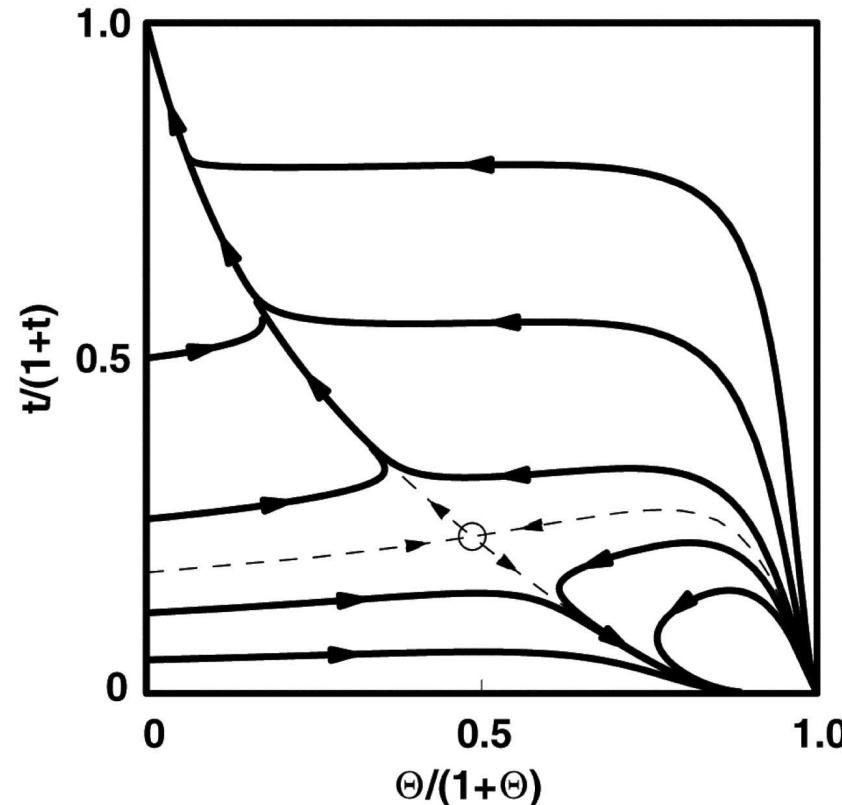
## Electron-electron interaction

E-e interaction can be incorporated  
within the same general theoretical framework ( $\sigma$  model);  
in some cases essentially modifies localization properties

# MIT in a 2D gas with strong interaction

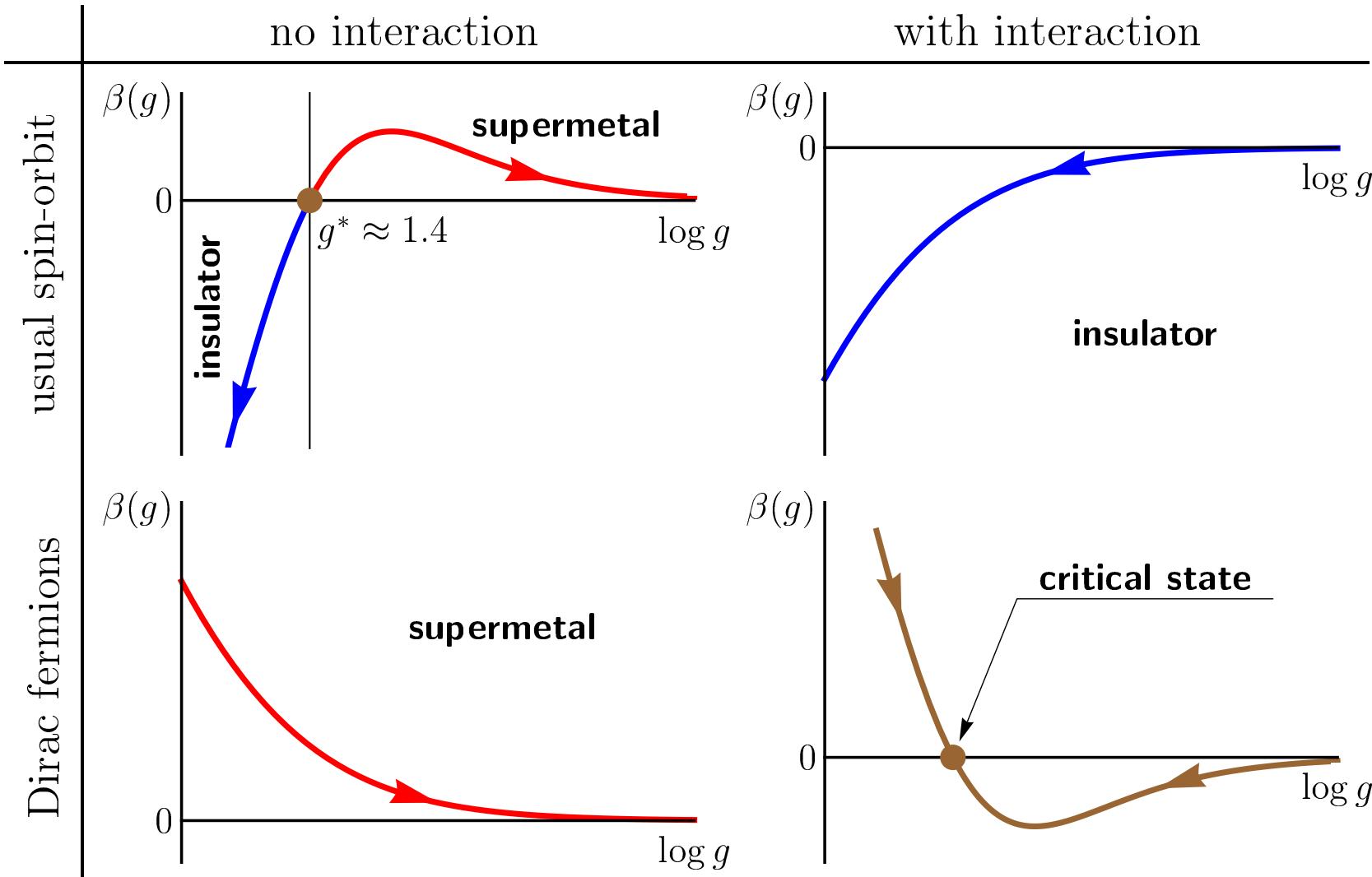


Kravchenko et al '94, ...

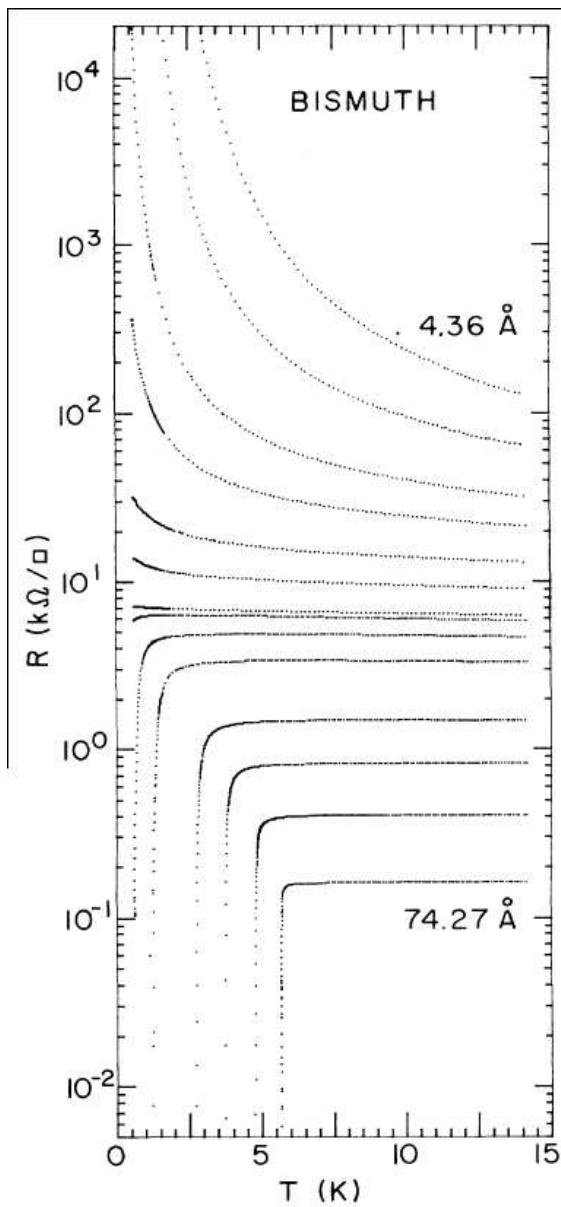


Finkelstein '83;  
Punnoose, Finkelstein, '02-05  
(number of valleys  $N \gg 1$ ;  
in practice,  $N = 2$  sufficient)

# Localization behavior in 2D Symplectic class: Effects of symmetry, topology, and Coulomb interaction

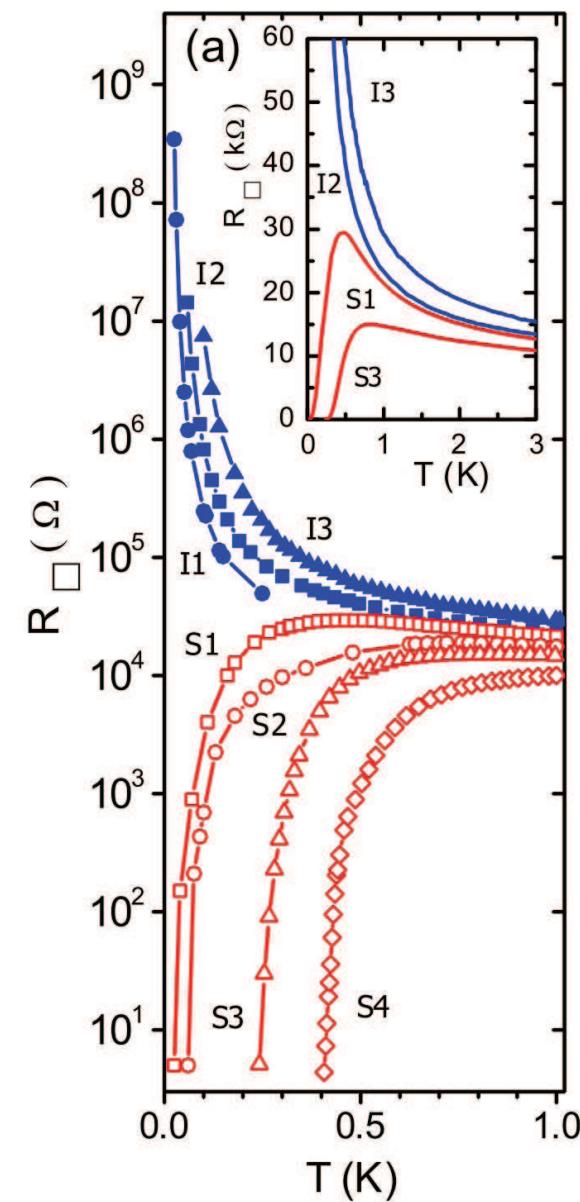
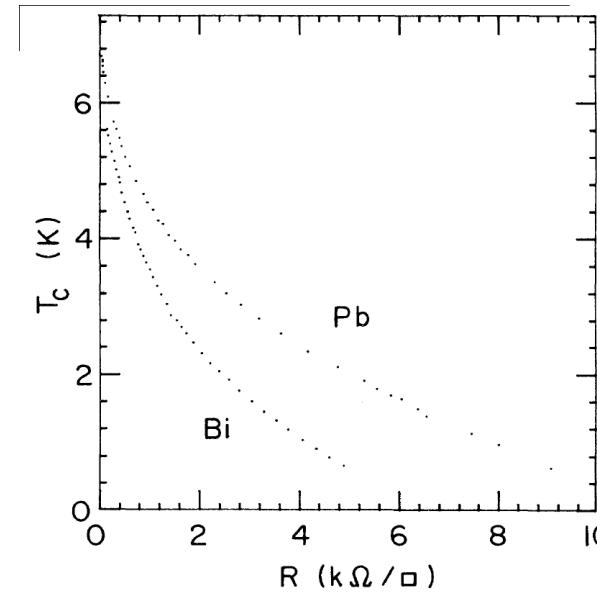


# Superconductor-Insulator Transition



Suppression of  $T_c$   
by disorder

Haviland, Liu, Goldman, PRL'89  
Bi and Pb films



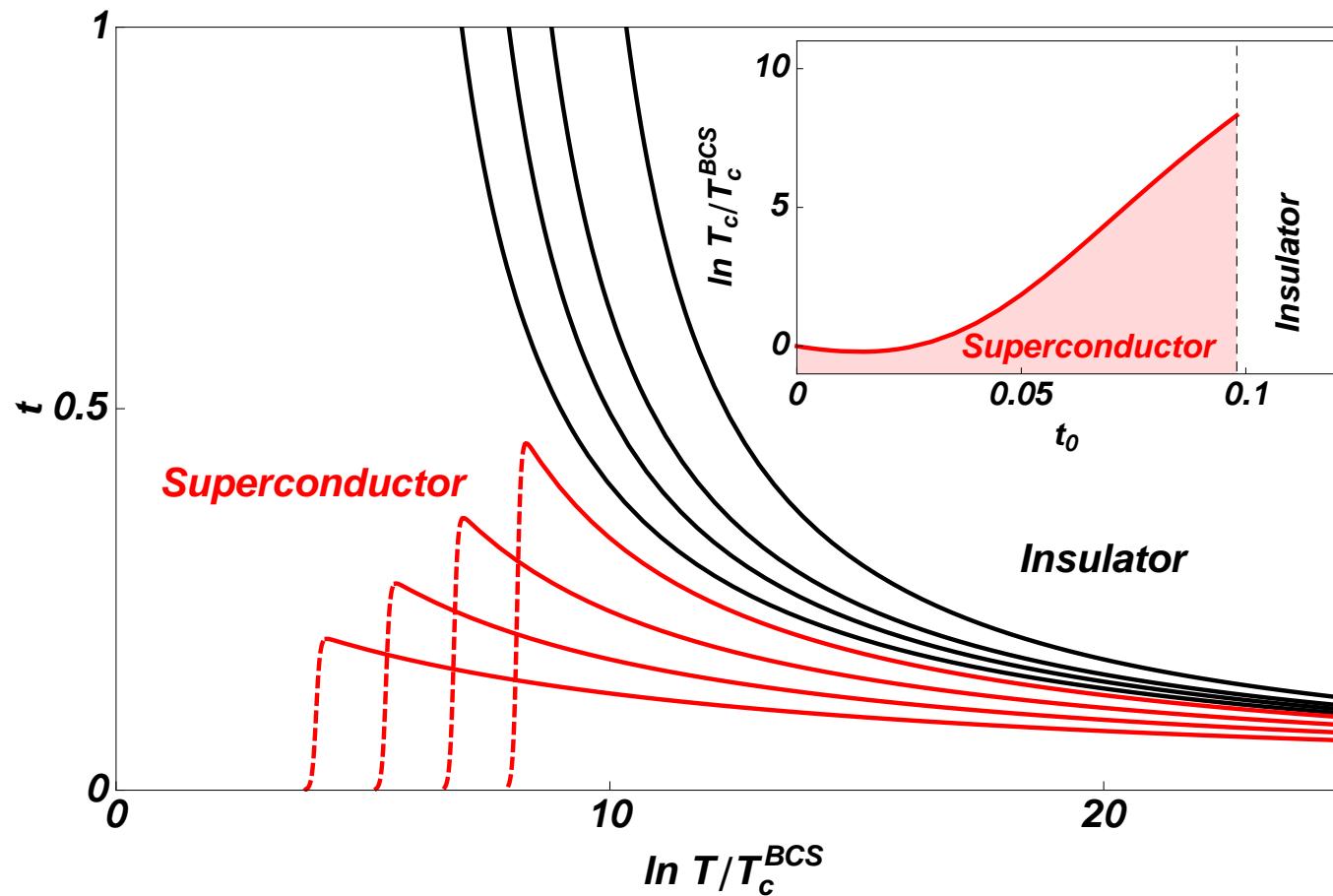
Baturina et al, PRL'07  
TiN films

# Enhancement of superconductivity by Anderson localization: interplay of interaction and multifractality

short-range interaction

Feigelman et al, PRL '07, Ann. Phys.'10: near 3D Anderson transition

Burmistrov, Gornyi, ADM, PRL '12: 2D films



# **International group on Localization, Interactions, and Superconductivity**

**Landau Institute for Theoretical Physics**

**RSF grant**

**“Superconductor-insulator and metal-insulator transitions  
in interacting disordered electronic systems” 2014-2016**

**Workshop 22-25 December 2014, Chernogolovka**

## Summary

- Anderson localization: basic properties, field theory
- Wave function multifractality
- Symmetries of disordered systems
- Manifestations of topology in localization theory
- Influence of electron-electron interaction

## Collaboration:

F. Evers, A. Mildenberger, I. Gornyi, P. Ostrovsky, I. Protopopov,  
E. König, S. Bera (Karlsruhe)  
M. Titov (Edinburgh – Nijmegen)  
S. Gattenlöhner, W.-R. Hannes (Edinburgh)  
M. Zirnbauer (Köln)  
I. Gruzberg, A. Subramaniam (Chicago)  
Y. Fyodorov (Nottingham – London)  
A. Ludwig (Santa Barbara)  
I. Burmistrov (Moscow)